Amplitudes in D=6 N=(1,1) SYM Theory

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Maximal SYM



- UV divergences First UV divergent diagrams cancel) Conformal or dual conformal symmetry Drummond, Hern, Korchemsky, Sokatchev Common structure of the integrande Bern, Dixon & Co 10

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BDS conjecture Bern, Dixon, Smirnov 05

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BDS conjecture

$$\mathcal{M}_{n} \equiv \frac{A_{n}}{A_{n}^{tree}} = 1 + \sum_{L=1}^{\infty} \left(\frac{g^{2} N_{c}}{16\pi^{2}}\right)^{L} M_{n}^{(L)}(\epsilon) = \exp\left[\sum_{l=1}^{\infty} \left(\frac{g^{2} N_{c}}{16\pi^{2}}\right)^{l} \left(f^{(l)}(\epsilon) M_{n}^{(1)}(l\epsilon) + C^{(l)} + E_{n}^{(l)}(\epsilon)\right)\right]$$

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$$\mathcal{M}_{n}(\epsilon) = \exp\left[-\frac{1}{8} \sum_{l=1}^{\infty} \left(\frac{g^{2} N_{c}}{16\pi^{2}} \right)^{l} \left(\frac{\gamma_{cusp}^{(l)}}{(l\epsilon)^{2}} + \frac{2G_{0}^{(l)}}{l\epsilon} \right) \sum_{i=1}^{n} \left(\frac{\mu^{2}}{-s_{i,i+1}} \right)^{l\epsilon} + \frac{1}{4} \sum_{l=1}^{\infty} \left(\frac{g^{2} N_{c}}{16\pi^{2}} \right)^{l} \gamma_{cusp}^{(l)} F_{n}^{(1)}(0) + C(g) \right]$$

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D=6
$$[g^2] \sim \frac{1}{M^2}$$

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 $[g^2] \sim \frac{1}{M^2}$ **D=6** Toy model for gravity

Color decomposition & Spinor helicity formalism

Color ordered amplitude

$$\mathcal{A}_n^{a_1\dots a_n}(p_1^{\lambda_1}\dots p_n^{\lambda_n}) = \sum_{\sigma \in S_n/Z_n} Tr[\sigma(T^{a_1}\dots T^{a_n})] \mathcal{A}_n(\sigma(p_1^{\lambda_1}\dots p_n^{\lambda_n})) + \mathcal{O}(1/N_c)$$

Planar Limit $N_c \to \infty, g_{YM}^2 \to 0 \text{ and } g_{YM}^2 N_c$ - fixed

Spinor helicity formalism

Cheung, O'Connell 09, Bern&Co 10

Momentum
$$p^{\mu}$$
, $p^2 = 0$, $\mu = 0, ..., 5$ $SO(5,1)$ $p_{AB} = p_{\mu}(\sigma^{\mu})_{AB}$, $p^{AB} = p^{\mu}(\bar{\sigma}_{\mu})^{AB}$ $SU(4)$ λ $p^{AB} = \lambda^{Aa}\lambda^B_a$, $p_{AB} = \tilde{\lambda}^{\dot{a}}_A \tilde{\lambda}_{B\dot{a}}$ Little group in D=6: $SO(4) \simeq SU(2) \times SU(2)$ Lorentz invariant structures: $\lambda(i)^{Aa} \tilde{\lambda}(j)^{\dot{a}}_A \doteq \langle i_a | j_{\dot{a}}] = [j_{\dot{a}} | i_a \rangle$

 $[1_{\dot{a}}2_{\dot{b}}3_{\dot{c}}4_{\dot{d}}] \doteq \epsilon^{ABCD}\tilde{\lambda}^{\dot{a}}_{A.1}\tilde{\lambda}^{\dot{b}}_{B.2}\tilde{\lambda}^{\dot{c}}_{C.3}\tilde{\lambda}^{\dot{d}}_{D.4}$

$$\langle 1_a 2_b 3_c 4_d \rangle \doteq \epsilon_{ABCD} \lambda_1^{Aa} \lambda_2^{Bb} \lambda_3^{Cc} \lambda_4^{Dd}$$

3

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Helicity is no longer conserved in D=6!

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Lorentz invariant structures:

$$\langle 1_a 2_b 3_c 4_d \rangle \doteq \epsilon_{ABCD} \lambda_1^{Aa} \lambda_2^{Bb} \lambda_3^{Cc} \lambda_4^{Dd}$$

SO(5, 1)

$$SU(4)$$
 λ^{Aa}
 $SO(4) \simeq SU(2) \times SU(2)$

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Superfield formalism in D=6



Universal expansion for any D in maximal SYM

Exact calculation

$$p_{i}^{2} = 0, \ m = 0$$

$$B_{1}(s,t) = \frac{\pi^{3}}{(2\pi)^{6}} \frac{b_{2}(x)}{s+t}, \ b_{2}(x) = \frac{L^{2}(x) + \pi^{2}}{2}, \ L(x) \doteq \log(x), \ x = \frac{t}{s}$$

$$B_{2}(s,t) = \left(\frac{\pi^{3}}{(2\pi)^{6}}\right)^{2} \left(\frac{b_{4}(x)}{t} + \frac{b_{3}(x)}{s+t}\right) \qquad \text{Anastasiou, Tausk, Tejeda-Yeomans, 00}$$

$$B_{2}(s,t) = \left(2\zeta_{3} - 2Li_{3}(-x) - \frac{\pi^{2}}{3}L(x)\right)L(1+x) + \left(\frac{1}{2}L(x) + \frac{\pi^{2}}{2}\right)L^{2}(1+x)$$

$$+ \left(2L(x)L(1+x) - \frac{\pi^{2}}{3}\right)Li_{2}(-x) + 2L(x)S_{1,2}(-x) - 2S_{2,2}(-x)$$

$$b_{3}(x) = -2\zeta_{3} + \frac{\pi^{2}}{3}L(x) - \left(L(x) + \pi^{2}\right)L(1+x) - 2L(x)Li_{2}(-x) + 2Li_{3}(-x)$$

Regge Limit $s \to \infty$, t < 0, fixed

$$B_1(s,t) \sim \frac{1}{2}L^2(x)$$
 $B_2(s,t) \sim \frac{1}{12}L^4(x)$

Leading Logarithms

UV finite

 $\frac{A_4}{A_4^{(0)}}$

Regge Limit $s \to \infty$, t < 0, fixed

$$B_n(t,s) \simeq rac{1}{s} rac{L^{2n}(x)}{n!(n+1)!}, \quad L \equiv \log(s/t)$$
 Bork, Kazakov, Vlasenko, 13

$$\frac{A_4}{A_4^{(0)}} \bigg|_{L.L.} = \sum_{n=0}^{\infty} \frac{(-g^2 t/2)^n L^{2n}(x)}{n!(n+1)!}, \quad \text{where} \quad g^2 \equiv \frac{g_{YM}^2 N_c}{64\pi^3}$$

$$\sum_{n=0}^{\infty} \frac{(-g^2 t/2)^n L^{2n}(x)}{n!(n+1)!} = \frac{I_1(2y)}{y}, \quad y \equiv \sqrt{g^2 |t|/2} \ L(x)$$

Regge behaviour

Exact for $N_c \to \infty$

$$\alpha(t) = 1 + 2\sqrt{g^2|t|/2} = 1 + \sqrt{\frac{g_{YM}^2 N_c|t|}{32\pi^3}}$$

 $\alpha(t)-1$

 $\sim \left(\frac{s}{-1}\right)$

$$B_n(s,t) = rac{1}{s} \left(C_n + O(t/s)
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 Kazakov, 14



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Leading Divergences

Loops	Combinatorics	Divergence
3	$(-g^2s/2)^3 \ 2t/s$	$1/6\epsilon$
4	$(-g^2s/2)^4 \ 2t/s$	$1/36\epsilon^2$
5	$(-g^2s/2)^5 \ 2t/s$	$1/216\epsilon^3$

Geom progression !?

$$\frac{A_4}{A_4^{(0)}} \bigg|_{Leading Div.} = 2\frac{t}{s} \sum_{n=1}^{\infty} \left(-\frac{g^2 s}{2}\right)^{n+2} \left(\frac{1}{6\epsilon}\right)^n = 2\frac{t}{s} \left(-\frac{g^2 s}{2}\right)^2 \frac{\frac{-g^2 s}{12\epsilon}}{1+\frac{g^2 s}{12\epsilon}}$$

 $\epsilon \to +0$

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In the limit ϵ ->0 the full expression is FINITE !

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That would imply that severe UV divergences present in any given order of PT are actually artifacts of the weak coupling expansion.

Figure 16 Figure 16 If this is true, one may try to apply the same arguments to quantum gravity.

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In order to understand the nonrenormalizable theories one has to find an alternative description.

The result of an alternative approach might be quite different from the PT one.