

Outline

Motivations

Classical scale  
invariance

Unimodular  
gravity

Phenomenology

Quantum  
scale  
invariance

Conclusions  
and Outlook

# A cosmological model based on scale invariance and unimodular gravity

Daniel Zenhäusern, EPF Lausanne

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Work in collaboration with Prof. Mikhail Shaposhnikov

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- 2 Classical scale invariance
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- 6 Conclusions and Outlook

## Problems in "standard" cosmology

- Horizon problem, flatness problem, relic problem,...  
→ *Inflation*
- Missing matter  
→ *Dark matter*
- Accelerated expansion of the universe  
→ *Dark energy*
- Cosmological constant problem (from qft)  
→ ?

There is enough motivation to modify something!

# Guiding idea

- In the SM and GR there are many different mass scales:
  - $m_{\text{Higgs}}, \langle h \rangle$
  - $M_P \sim \frac{1}{\sqrt{G}}$
  - $\Lambda_{\text{QCD}}, \dots$
- Question 1: Can we construct a theory in which all these scales are related, induced dynamically?
- Question 2: Is it possible to construct a theory that doesn't contain such absolute scales at all?

## Construction of the model

Introduce a new scalar field  $\chi$ .

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{SM}[M \rightarrow 0]} + \mathcal{L}_G + \frac{1}{2}(\partial_\mu \chi)^2 - V(\varphi, \chi) \\ + (\bar{N}_I i \gamma^\mu \partial_\mu N_I - h_{\alpha I} \bar{L}_\alpha N_I \tilde{\varphi} - f_I \bar{N}_I^c N_I \chi + \text{h.c.}) ,$$

$$V(\varphi, \chi) = \lambda \left( \varphi^\dagger \varphi - \frac{1}{2} \zeta^2 \chi^2 \right)^2 + \beta (\chi^2 - \chi_0^2)^2 ,$$

$$\mathcal{L}_G = - \left( \xi_\chi \chi^2 + 2\xi_h \varphi^\dagger \varphi \right) \frac{R}{2} .$$

$\chi_0 \neq 0, \beta > 0 \Rightarrow$  Induced gravity (Zee (1979), Smolin (1979))

# The case for scale invariance

- The model is somewhat artificial.
  - Terms  $m_H^2 \varphi^\dagger \varphi$ ,  $m_N \bar{N}_I^c N_I$  and  $M_P^2$  are absent.
  - Terms  $\chi_0^2 \chi^2$  and  $\chi_0^4$  are present.
- Try to construct a theory with no dimensional constants at all  $\rightarrow \chi_0 = 0$ .
- No dimensional parameters  $\rightarrow$  dilatational symmetry

$$\phi(x) \rightarrow \sigma^n \phi(\sigma x)$$

$$g_{\mu\nu}(x) \rightarrow g_{\mu\nu}(\sigma x)$$

## The case for a flat direction

$$V(\varphi, \chi) = \frac{\lambda}{4} (h^2 - \zeta^2 \chi^2)^2 + \beta \chi^4$$

- The ground state should break the dilatational symmetry.
- "Only possibility" is to have a flat direction  $\rightarrow \beta = 0$ .

$$V(\varphi, \chi) = \frac{\lambda}{4} (h^2 - \zeta^2 \chi^2)^2$$

- Ground state is infinitely degenerate.
- Physics doesn't depend on field values at the minimum.

## Interim balance

Model based on scale invariance + flat direction

- Horizon problem, flatness problem, monopole problem,...  
→ *Inflation due to roll down to the potential minimum*
- Missing matter  
→ *Dark matter given by sterile neutrinos*
- Cosmological constant problem  
→  $\Lambda = 0$ , *due to scale invariance and flat direction*
- Accelerated expansion of the universe  
→ *Problem is not adressed by this model!*

**We need another ingredient.**



# Unimodular gravity

- UG mainly appeared in the context of the cosmological constant problem

$$\mathcal{L} = \sqrt{-g}(-M_P^2 R + \Lambda)$$

- Idea is to constrain variables by  $\det(g_{\mu\nu}) = -1$

$$\mathcal{L} \rightarrow \mathcal{L}_{UG} = -M_P^2 \hat{R} + \Lambda$$

- Equations of motion are

$$\hat{R}_{\mu\nu} - \frac{1}{4}\hat{R}\hat{g}_{\mu\nu} = 0$$

→  $\Lambda$  *doesn't appear.*

- But apply  $\hat{\nabla}^\mu$  to equations  $\Rightarrow \partial_\nu \left( \frac{\hat{R}}{4} \right) = 0$   
 $\Rightarrow \frac{\hat{R}}{4} = \text{constant} \equiv \Lambda_0$
- Reinsert this into equations to get

$$\hat{R}_{\mu\nu} - \frac{1}{2}\hat{R}\hat{g}_{\mu\nu} + \Lambda_0\hat{g}_{\mu\nu} = 0$$

$\rightarrow$  *Integration constant acts as a cosmological constant.*

- These are the equations for GR with a cosmological constant  $\Lambda_0$  in the coordinate frame where  $\det g_{\mu\nu} = -1$ .
- Hence, the result

$$\hat{\mathcal{L}} = -M_P^2 \hat{R} \quad \leftrightarrow \quad \mathcal{L} = \sqrt{-g}(-M_P^2 R + \Lambda_0)$$

(Van der Bij (1982), Zee (1985),...)

# UG plus other fields

- General result for UG in combination with non-gravitational fields.

$$\mathcal{L}_{UG} = \mathcal{L}(\hat{g}_{\mu\nu}, \phi, \partial\phi) \leftrightarrow \mathcal{L}_{GR} = \sqrt{-g}(\mathcal{L}(g_{\mu\nu}, \phi, \partial\phi) + \Lambda_0)$$

Result doesn't depend on the nature of the fields and on the way they couple to gravity.

## UG in our model

- Replace GR by UG, scalar and gravity parts become

$$\mathcal{L}_{UG} = -\frac{1}{2}(\xi_\chi \chi^2 + \xi_h h^2) \hat{R} + \frac{1}{2}(\partial\chi)^2 + \frac{1}{2}(\partial h)^2 - V(h, \chi)$$

- Theory is still scale invariant.
- Has the same classical solutions as

$$\mathcal{L}_{GR} = \sqrt{-g} \left( -\frac{1}{2}(\xi_\chi \chi^2 + \xi_h h^2) R + \frac{1}{2}(\partial\chi)^2 + \frac{1}{2}(\partial h)^2 - V(h, \chi) - \Lambda_0 \right)$$

$\Lambda_0$  depends on initial conditions and is not really part of the action.

- Classically it's enough to analyse  $\mathcal{L}_{GR}$  for different values of  $\Lambda_0$ .

# UG in our model

- Change to the Einstein frame

$$\tilde{g}_{\mu\nu} = (\xi_\chi \chi^2 + \xi_h h^2) M_P^{-2} g_{\mu\nu}$$

- Action in Einstein frame

$$\mathcal{L}_E = \sqrt{-\tilde{g}} \left( -M_P^2 \frac{\tilde{R}}{2} + K - U_E(h, \chi) \right)$$

$$U_E(h, \chi) = \frac{M_P^4}{(\xi_\chi \chi^2 + \xi_h h^2)^2} (V(h, \chi) + \Lambda_0)$$

$K$  = kinetic term

- $K$  is non diagonal, but positive definite for  $\xi_\chi, \xi_h > 0$ .

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Classical scale invariance

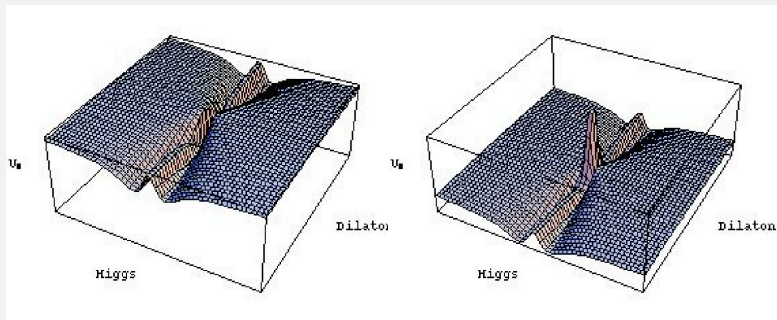
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$$U_E(h, \chi) = \frac{M_P^4}{(\xi_\chi \chi^2 + \xi_h h^2)^2} \left( \frac{\lambda}{4} (h^2 - \zeta^2 \chi^2)^2 + \Lambda_0 \right)$$



$$\Lambda_0 < 0$$

$$\Lambda_0 > 0$$

We will suppose that  $0 < \Lambda_0 \lesssim M_P^4$

# Inflation and reheating

- $$U_E(h, \chi) = \frac{M_P^4}{(\xi_\chi \chi^2 + \xi_h h^2)^2} (V(h, \chi) + \Lambda_0)$$
- Chaotic inflation  $\rightarrow$  initially  $\chi$  and  $h$  away from valley and  $\chi, h > M_P \rightarrow$  First term in potential dominates.
- For  $\xi_h \gg 1, \xi_\chi \ll 1$  and  $\chi \sim h$  dynamics dominated by  $h \rightarrow \chi \sim \text{const.}$
- Inflation due to roll down of  $h$  into the valley.
- Reheating due to oscillations of  $h$  in the valley.
- $\xi_h \sim 700 - 20000$  produces correct perturbation spectrum and a reheating temperature of  $T_{rh} \sim 10^{13} \text{ GeV.}$   
(Bezrukov and Shaposhnikov)

## Late time evolution

- Fields are trapped in the valley  $h \sim \zeta \chi$ .  
→ Second term of potential dominates.
- $U_E(h, \chi) \sim M_P^4 \frac{\Lambda_0}{(\xi_\chi \chi^2 + \xi_h h^2)^2}$
- Good phenomenology can be obtained for  
 $\lambda \sim 1$  and  $\zeta \lll 1 \Rightarrow h \lll \chi$ .
- Define dilaton field  $\eta$  along the valley and  $\phi$  perpendicular to it.
- Equations of motion

$$\ddot{\eta} + 3H\dot{\eta} + \frac{dU_\eta}{d\eta} \approx 0$$

$$\ddot{\phi} + 3H\dot{\phi} + m_h^2 \phi \approx 0$$



## Late time evolution

The late time evolution of the universe is described by

$$\ddot{\eta} + 3H\dot{\eta} + \frac{dU_{\eta}}{d\eta} = 0$$

$$H^2 = \frac{1}{3M_P^2} \left( \frac{1}{2}\dot{\eta}^2 + U_{\eta} + \frac{C_{\gamma}}{a^4} + \frac{C_M}{a^3} \right)$$

$$U_{\eta} = \frac{\Lambda_0}{\xi_X^2} \exp\left(-\frac{\gamma\eta}{M_P}\right)$$

$$\gamma = \frac{4}{\sqrt{6 + \frac{1}{\xi_X}}}$$

→ well-studied model (quintessence field,...)

- $\gamma > \sqrt{3} \rightarrow$  scaling solutions, no dark energy
- $\gamma < \sqrt{3} \rightarrow$  no scaling solutions,  $\Omega_{\eta}$  keeps growing (Wetterich (1988), Ferreira and Joyce (1998))

## Dark energy

- For  $\gamma < \sqrt{3}$  thawing quintessence (Ferreira and Joyce (1998), Caldwell(2005))  $\rightarrow$  dark energy
- Field frozen by Hubble friction at early times, then slow roll-down
- Result for this scenario (Ferreira and Joyce (1998))

$$1 + \omega = \frac{\gamma^2}{3} \left[ \frac{1}{\sqrt{\Omega_\eta}} - \frac{1}{2} \left( \frac{1}{\Omega_\eta} - 1 \right) \log \frac{1 + \sqrt{\Omega_\eta}}{1 - \sqrt{\Omega_\eta}} \right]^2$$

- Identify  $\Omega_\eta = \Omega_{DE}$
- WMAP data  $\rightarrow -0.04 < 1 + \omega < 0.2$  and  $\Omega_{DE} \approx 0.73$   
 $\Rightarrow \gamma < 1.11 < \sqrt{3}$  (ok for thawing scenario)
- $0 < \xi_\chi < 0.16$
- Fate of the universe in this model  
 $\Omega_\eta \rightarrow 1, \quad \omega \rightarrow \frac{\gamma^2}{3} - 1, \quad a(t) \propto t^{\frac{2}{\gamma^2}}$   
 (Ferreira, Joyce)

# Interim balance

Model based on scale invariance + flat direction + **UG**

- Horizon problem, flatness problem, relic problem,...  
→ *Inflation due to roll down to the potential minimum*
- Missing matter  
→ *Dark matter given by sterile neutrinos*
- Cosmological constant problem  
→  $\Lambda = 0$  *since we require a flat direction*
- Accelerated expansion of the universe  
→ **Dark energy due to UG with  $\Lambda_0 > 0$**

## Quantum aspects

- Results remain valid on quantum level if
  - Quantum theory is exactly scale invariant
  - Dilaton remains exactly massless (flat direction)
  - Initial conditions of UG give dark energy
- But, in all common theories dilatation symmetry is anomalous
  - $\rightarrow \partial_\mu J^\mu \propto \beta(g) G_{\alpha\beta}^a G^{\alpha\beta a}$
- Can we construct a scale invariant quantum theory?

# Standard renormalization

- Theory is continued to  $d = 4 - 2\epsilon$  dimensions

$$S = \int dx^d \left\{ \frac{1}{2} [(\partial_\mu \chi)^2 + (\partial_\mu h)^2] - \lambda (h^2 - \zeta^2 \chi^2)^2 \right\}$$

- Fields have mass dimension  $1 - \epsilon$
  - $\lambda$  has dimension  $2\epsilon$
  - Action is no longer scale invariant.
- Define finite dimensionless coupling  $\lambda_R$

$$\lambda = \mu^{2\epsilon} \left[ \lambda_R + \sum_{n=0}^{\infty} \frac{a_n}{\epsilon^n} \right] \quad \mu \text{ has dimensions of mass}$$

- Fix  $a_n$  such that renormalized Green's functions are finite to all orders of perturbation theory.

# Standard renormalization

- One loop effective potential along the flat direction in  $\overline{\text{MS}}$  scheme

$$V_1(\chi) = \frac{m_H^4(\chi)}{64\pi^2} \left[ \log \frac{m_H^2(\chi)}{\mu^2} - \frac{3}{2} \right]$$

$$m_H^2(\chi) = 2\lambda\zeta^2(1 + \zeta^2)\chi^2$$

- Dilatation anomaly appears because regularization breaks the symmetry.

# Scale invariant procedure

- Continuation to d-dimensions in a scale invariant way.
- Instead of  $\mu$  introduce a combination of the fields with the correct mass dimension

$$\mu^{2\epsilon} \rightarrow (\xi_\chi \chi^2 + \xi_h h^2)^{\frac{\epsilon}{1-\epsilon}}$$

# One-loop effective potential

- Denote  $\omega^2 = \xi_\chi \chi^2 + \xi_h h^2$ .
- Write the classical potential in d dimensions

$$U = \frac{\lambda_R}{4} [\omega^2]^{\frac{\epsilon}{1-\epsilon}} [h^2 - \zeta_R^2 \chi^2]^2$$

- Introduce counter terms

$$U_{cc} = [\omega^2]^{\frac{\epsilon}{1-\epsilon}} \left[ Ah^2 \chi^2 \left( \frac{1}{\bar{\epsilon}} + a \right) + B\chi^4 \left( \frac{1}{\bar{\epsilon}} + b \right) + Ch^4 \left( \frac{1}{\bar{\epsilon}} + c \right) \right]$$

- We have parameters  $A, B, C, a, b, c$ .



# One-loop effective potential

- $A, B, C$  are chosen to cancel the divergences.
- One loop correction is automatically scale invariant.
- $U_1(h, \chi)$  does not automatically have a flat direction.
- Fix  $b, c$  such that  $W_1(\zeta_R) = W_1'(\zeta_R) = 0$ .  
→ Flat direction

## Effective Higgs potential

One loop potential for  $\chi = \chi_0$ ,  $v \equiv \zeta_R \chi_0$  and  $\zeta_R \lll 1$

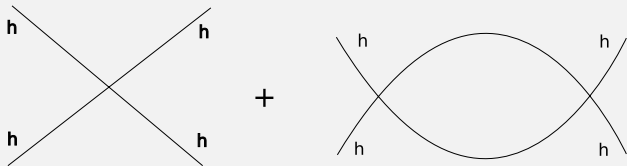
$$U_1 = \frac{m^4(h)}{64\pi^2} \left[ \log \frac{m^2(h)}{v^2} + \mathcal{O}(\zeta_R^2) \right] \\ + \frac{\lambda_R^2}{64\pi^2} [C_0 v^4 + C_2 v^2 h^2 + C_4 h^4] + \mathcal{O}\left(\frac{h^6}{\chi_0^2}\right)$$

where  $m^2(h) = \lambda_R(3h^2 - v^2)$ .

- First term is standard effective potential for the theory with  $h$  only and the potential  $U = \frac{\lambda}{4} (h^2 - \zeta^2 \chi_0^2)^2$ .
- Additional terms give redefinitions of coupling, mass and vacuum energy.
- Corrections to the Higgs mass proportional to  $v^2 \propto \zeta_R^2 \chi_0^2$   
→ No problem of stability of the Higgs mass.

# Running coupling

- Higgs-Higgs scattering at  $v \ll \sqrt{s} \ll \chi_0$  ( $\zeta_R \lll 1$ )



$$\Gamma_4 = \lambda_R + \frac{9\lambda_R^2}{64\pi^2} \left[ \log \left( \frac{s}{\xi_X \chi_0^2} \right) + \text{const} \right] + \mathcal{O}(\zeta_R^2)$$

- Coupling runs in standard way.

# Conclusions

- Constructed a scale invariant cosmological model.
  - All Scales are generated dynamically.
  - Inflation similar to Higgs inflation.
  - Sterile Neutrino as DM candidate.
  - Dark energy due to unimodular gravity.
  - $\omega > -1$
- Constructed a class of scale invariant quantum theories.
  - Dilatation symmetry is spontaneously broken.
  - Standard running of coupling constants.
  - No problem of stability of the Higgs mass.
- Classical results could remain valid in full quantum theory.

## Some remaining tasks and questions

- Detailed analysis of inflation.
- The model doesn't address the question about the big difference between the electroweak and the Planck scale.
- Can one do an experiment to detect effects of the dilaton?
- Are the scale invariant theories unitary, ... ?

**Thank You.**