Two – component – liquid model for quark – gluon plasma

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Model, reminder

In hydrodynamic approximation:

$$j^{\mu} = nu^{\mu} + f^{2}\partial^{\mu}\phi \qquad (1)$$
$$T^{\mu\nu} = (\epsilon + P)u^{\mu}u^{\nu} + P\eta^{\mu\nu} + f^{2}\partial^{\mu}\phi\partial^{\nu}\phi \qquad u^{\mu}\partial_{\mu}\phi = \mu$$

 $\begin{array}{l} \pmb{u}^{\mu} \text{ is 4-velocity,} \\ \phi \text{ is a scalar field} \end{array}$

Two independent motions,

$$\rho_{tot} = \rho_n + \rho_s$$

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Outline of the talk

- **1** Why at all?
- **2** where is the scalar field in QCD?

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3 possible crucial test

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Basic properties of plasma

A. Equation of state

$$\epsilon(T) \approx (\epsilon(T))_{ideal gas} (1 - \delta)$$

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 $\delta \approx 0.1 - 0.15$

Close to the ideal gas

Known since long from lattice simulations

"lost decades" (E.V. Shuryak)

Two — component — liquid model for quark — gluon plasma Basic properties of plasma

Viscosity η

(like friction, enters hydrodynamic $T^{\mu\nu}$) Fits to the RHIC data:

$$\left(rac{\eta}{s}
ight)_{\it plasma}~pprox~rac{1}{4\pi}$$

It is the lowest viscosity among all known liquids (or, the plasma is closest to the ideal liquid)

For ideal gas:

$$\eta_{\it ideal\ gas} \
ightarrow \ \infty$$

Two — component — liquid model for quark — gluon plasma └─Basic properties of plasma

C. Quantum effects

From kinetics:

$$\frac{\eta}{\mathbf{s}} \sim \mathbf{k}_{B}^{-1} \tau_{relaxation} \cdot \left(\frac{\epsilon}{n}\right)$$

From uncertainty principle

$$\left(rac{\epsilon}{n}
ight)\cdot au_{ ext{relaxation}} ~\sim~ ar{h}$$

since ϵ/n is energy per particle

$$rac{\eta}{m{s}} \sim rac{ au$$
relaxation}{ au_{quantum}}

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It is a challenge to explain points A)-C) which show in opposite directions

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Two components?

What is special about the viscosity? If there are two components

$$\frac{1}{\eta_{tot}} = \frac{c_1}{\eta_1} + \frac{c_2}{\eta_2} \qquad (*)$$

where $c_1 + c_2 = 1$, $c_{1,2}$ are phase-space factors Indeed, η is like resistance

(*) is true for classical solutions (*c*_{1,2} are concentrations)
(*) is supported by superfluidity example

Two components:

- **1** One component dominates E.o.S., $c_1 >> c_2$
- **2** The other dominates η_{tot} if $\eta_2 \ll \eta_1$
- **3** if one component superfluid, large quantum effects are 'explained'
- unlike non-relativistic case, in field th. seems no limit on temperature

Indeed, from dimensional reduction:

$$\lim_{T\to\infty} c_2 ~\sim~ \frac{1}{(\ln T)^3}$$

To summarize: two components seem to work qualitatively

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Scalar field?

To realize superfluidity, need a scalar fields. Constraints:

- **1** complex field,
 - $\phi \ \neq \phi^* \ ,$
- 2 condensed,

$$\langle \phi \rangle \neq \mathbf{0} ,$$

3 no known quantum number is allowed to be violated by $\langle \phi \rangle \neq \mathbf{0}$,

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4 rather **3***d* field, $\phi(\mathbf{r})$ ($\partial_t \phi = \mu$)

Thermal scalar

Scalars do arise within string-based approaches Near Hagedorn transition, $\beta < \beta_H$ single mode with mass

$$m_{eta}^2 = rac{eta_{H}(eta_{H}-eta)}{2\pi^2(lpha')^2}$$

dominates. At $\beta = \beta_H$ becomes tachyonic.

$$F = = \beta \ln Z \approx \beta \int_0^\infty \frac{dL}{L} \frac{\exp(-m_\beta^2 l_s L)}{(l_s L)^{d/2}}$$

Sum over random walks

- 1 of length L
- **2** with steps I_s (related to the string tension)
- 3 in **d** dimensions

Two – component – liquid model for quark – gluon plasma └─Scalar fields from strings

Thermal scalar vs constraints

IF thermal scalar condensed at $T > T_c$ our constraints were satisfied:

- **2** 3d field (from explicit calculations, strings static)
- quantum number related to the field is topological (wrapping around the time direction)

Gauge/string correspondence

Duals to YM: strings in extra curved dimensions

Generic features:

- **I** Thermal scalar condensation changes geometry
- 2 various lower-dimensions defects exist
- **3** at $T > T_c$ defects become time oriented (strings static)
- 3d projections of time-oriented strings do produce scalar fields
- **5** Deconfinement phase transition as a change of geometry from 4d to 3d already at $T = T_c$

Two – component – liquid model for quark – gluon plasma └─Scalar fields from strings

Lattice data

provide independent support for theoretical picture above

In particular, 2d defects become time oriented at $T>T_c$ (known since long)

At a time slice their projection forms an infinite cluster that is,

scalar condensation is observed in geometric language

└─Scalar fields from strings

To summarize:

there is strong evidence in favor of existence of 3d scalar field at $T > T_c$

both from continuum (dual models) and lattice.

Deconfinement phase transition as transition from 4d to 3d in non-perturbative sector.

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Two - component - liquid model for quark - gluon plasma └ Possible crucial test

Static correlator of momentum densities operators

Consider

$$G_{R}^{0j,0i}(k) \;=\; i\int d^{4}x e^{-ikx} heta(t) < |T^{0j}(x),T^{0i}(0)| >$$

in the static case:

$$G_R^{0j,0i}(\omega=0,\mathbf{k}) = rac{k_ik_j}{\mathbf{k}^2}G^L(\mathbf{k}) + \left(\delta^{ij} - rac{k_ik_j}{\mathbf{k}^2}
ight)G^T(\mathbf{k})$$

Theorems known:

$$\lim_{\mathbf{k}\to \mathbf{0}} \boldsymbol{G}^{\mathsf{T}}(\mathbf{k}) = -(\boldsymbol{s}\boldsymbol{T} + \mu\rho_{\mathsf{n}})$$

and

$$\lim_{L \to \infty} G^{L}(\mathbf{k}) = -(sT + \mu \rho_{tot}) \oplus (s + s) \oplus$$

Possible crucial test

In other words,

$$\lim_{\mathbf{k}=0} G^{0j,0i}(\mathbf{k}=0) = \rho_s \frac{k_i k_j}{\mathbf{k}^2}$$

Superfluidity brings in non-analyticity at small \mathbf{k} . All quantities are static and could, therefore, be measured directly on the lattice

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L_{Conclusions}

The two-component model has not failed so far.

Could be crucially tested in the future