### 5d Models of Mesons and Baryons

Andrea Wulzer

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#### Outline

Introduction A 5d model of mesons Skyrmions in 4 and 5d Conclusions

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A 5d model of mesons

Skyrmions in 4 and 5d

Conclusions

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Can a weakly-coupled (dual) description of the hadrons exist?

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Yes, but only in the large-N<sub>c</sub> expansion: ('t Hooft, Veneziano, Coleman, Witten ...)

- Meson couplings scale like  $1/\sqrt{N_c}$
- Meson masses scale like  $N_c^0$
- Mesons come in infinite towers

Weakly interacting theory of an  $\infty$  number of mesons.

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Weakly interacting theory of an  $\infty$  number of mesons.

Weak coupling:  $g_{
ho} \simeq 4\pi/\sqrt{N_c} \rightarrow 0$ 

 $1/N_c$  is the only known candidate coupling for a dual theory

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Extra-dimensions are the only tool we have to describe towers. Do simple field theories in 5d capture some features of large- $N_c$ ?

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Attractiveness of the 5d models:

- 1. "Holographic" implementation of  $\chi_{SB}$  automatically leads to KK towers of vector (and scalar) mesons.
- 2. Extremely predictive framework (description of  $\rho$ ,  $\omega$ ,  $a_1$ ,  $\rho'$  ...)
- 3. Easy bookkeping of  $1/N_c$  factors:

 $1/\sqrt{N_c} \sim 5 d$  coupling

4. 5d models are valid effective theories:

 $\Lambda_5/m_
ho 
ightarrow \infty$  for  $N_c 
ightarrow \infty$ 

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## Main Limitation:

Absence of high spin states and of Regge phenomenology Higher spins above the 5d cutoff, no known  $large-N_c$  reason String models could do better, but we have no candidate

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## Where are the Baryons?

Baryons are solitons (skyrmions) at large– $N_c$ (Witten)

- size ~  $N_c^0/\Lambda_{QCD}$
- ► mass ~ N<sub>c</sub> Λ<sub>QCD</sub>

Are there skyrmions in the 5d model?

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$$Z_{QCD} = \int \mathcal{D} \mathbf{L}_M \mathcal{D} \mathbf{R}_M e^{iS_5[\mathbf{L},\mathbf{R}]} \qquad \begin{matrix} UV & U(2)_L \times U(2)_R & IR \\ \downarrow & \downarrow & \downarrow \\ L = I & U(2)_V \\ R = r & L = R \end{matrix}$$

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The 5d action is:  $S_5 = S_{CS} + S_g$ 

$$\begin{split} S_{CS} &= -i \frac{N_c}{24\pi^2} \int \left[ \omega_5(\mathbf{L}) - \omega_5(\mathbf{R}) \right] \,, \quad \text{reproduces the QCD Anomaly} \\ S_g &= -\int \frac{M_5}{2} \left\{ \text{Tr} \left[ L_{MN} L^{MN} \right] \,+ \, \frac{\alpha^2}{2} \widehat{L}_{MN} \widehat{L}^{MN} \,+ \, \{L \,\leftrightarrow \, R\} \right\} + \dots \,. \end{split}$$

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 $S_{CS} = -i \frac{N_c}{24\pi^2} \int [\omega_5(\mathbf{L}) - \omega_5(\mathbf{R})], \text{ reproduces the QCD Anomaly}$   $S_g = -\int \frac{M_5}{2} \left\{ \text{Tr} \left[ L_{MN} L^{MN} \right] + \frac{\alpha^2}{2} \widehat{L}_{MN} \widehat{L}^{MN} + \{L \leftrightarrow R\} \right\} + \dots.$ Metric:  $ds^2 = a(z)^2 (\eta_{\mu\nu} dx^{\mu} dx^{\nu} - dz^2), a(z) = \frac{L}{z}, z \in [0, L]$ At two derivatives order we have 3 parameters:  $M_5, L, \alpha$ .

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#### Some interesting features:

Vector Meson Dominance holds in our model.

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Some interesting features:

- Vector Meson Dominance holds in our model.
- Decay constants and meson couplings all scale like:

 $F_i \sim \sqrt{M_5}, \quad g_i \sim 1/\sqrt{M_5}$ 

Correct  $N_c$  scalings if  $\alpha, L \sim N_c^0$  and  $M_5 \sim N_c$ .

- Automatic Zweig Rule for  $m_{\omega} = m_{\rho}$  only,  $F_{\omega} = \alpha F_{\rho}$ .
- "Modified" KSFR relation:  $m_{\rho}^2 \simeq 3g_{\rho\pi\pi}^2 F_{\pi}^2$  (exp=2.1).

Relations valid at the leading  $1/M_5~(1/N_c)$  order.

Though  $N_c = 3$ , what if we compare with real hadrons?

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## Comparison with experiments:

	Experiment	$\mathrm{AdS}_5$	Deviation
$m_{ ho}$	775	824	+6%
$m_{a_1}$	1230	1347	+10%
$m_\omega$	782	824	+5%
$F_{ ho}$	153	169	+11%
$F_{\omega}/F_{\rho}$	0.88	0.94	+7%
$F_{\pi}$	87	88	+1%
$g_{\rho\pi\pi}$	6.0	5.4	-10%
L <sub>9</sub>	$6.9 \cdot 10^{-3}$	$6.2 \cdot 10^{-3}$	-10%
L <sub>10</sub>	$-5.2 \cdot 10^{-3}$	$-6.2 \cdot 10^{-3}$	-12%
$\Gamma(\omega  ightarrow \pi \gamma)$	0.75	0.81	+8%
$\Gamma(\omega  ightarrow 3\pi)$	7.5	6.7	-11%
$\Gamma( ho  o \pi \gamma)$	0.068	0.077	+13%
$\Gamma(\omega  ightarrow \pi \mu \mu)$	$8.2 \cdot 10^{-4}$	$7.3\cdot10^{-4}$	-10%
$\Gamma(\omega  ightarrow \pi ee)$	$6.5 \cdot 10^{-3}$	$7.3 \cdot 10^{-3}$	+12%

Fit with 3 parameters to 14 observables, Total RMS Error of 11%. Observables selected to have < 10% experimental error.

## Skyrmions in 4 and 5d

## The Skyrme's Idea

Static Goldstones:  $U(\mathbf{x}): S^3 \to SU(2) \sim S^3, \quad \Pi_3(S^3) = Z$  $B = \frac{1}{24\pi^2} \int d^3x \, \epsilon^{ijk} \mathrm{Tr} \left[ U \partial_i U^{\dagger} \, U \partial_j U^{\dagger} \, U \partial_k U^{\dagger} \right] \in Z.$ 

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## Skyrmions in 4 and 5d

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B is exactly conserved topological charge:



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### Witten: Quantized Skyrmions have Spin/Isospin of the Baryons !!

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However, Skyrme's idea cannot be implemented in the  $\sigma$ -model:  $S_2 = \int \frac{F_{\pi}^2}{4} \operatorname{Tr} \left\{ \partial_{\mu} U \partial^{\mu} U^{\dagger} \right\} \rightarrow E_2[\rho] \propto F_{\pi}^2 \rho$  $S_4 \supset \int \frac{F_{\pi}^2}{\Lambda^2} \operatorname{Tr} \left\{ \left[ \partial_{\mu} U U^{\dagger}, \partial_{\nu} U U^{\dagger} \right]^2 \right\} \rightarrow E_4[\rho] \propto F_{\pi}^2 / \Lambda^2 1 / \rho$ 

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Skyrmion's size is  $\rho_S \propto 1/\Lambda$ . All operators contribute the same. Skyrmions are UV-dominated, crucially depend on resonances

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### The 5d Skyrmion (with A. Pomarol and G. Panico)

Topological charge:  $B = \frac{1}{32\pi^2} \int_0^L d^4 x \, \epsilon_{\hat{\mu}\hat{\nu}\hat{\rho}\hat{\sigma}} \operatorname{Tr} \left[ L^{\hat{\mu}\hat{\nu}} L^{\hat{\rho}\hat{\sigma}} - R^{\hat{\mu}\hat{\nu}} R^{\hat{\rho}\hat{\sigma}} \right]$ 

Can be shown equivalent to Skyrme's.

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Can be shown equivalent to Skyrme's.

The energy is: 
$$E = M_5 \int_0^L d^4 x \left[ F^2 + \gamma L \epsilon^{...} \widehat{A} FF \right]$$

For  $\gamma \ll 1$  expand around an instanton of size  $\rho$ :  $E(\rho) = 8\pi^2 M_5 \left[ 1 + \frac{\rho}{2L} + \gamma^2 \frac{L^2}{\rho^2} \right]$ 

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### The Chern–Simons term stabilizes the skyrmion: $\rho_S \simeq \gamma^{2/3} L$

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The Chern–Simons term stabilizes the skyrmion:  $\rho_S \simeq \gamma^{2/3} L$ Actually,  $\gamma = \frac{N_c}{64\pi^2 M_5 L} = \mathcal{O}(N_c^0)$  is not small Solution must be found numerically and  $\rho_S \simeq L \simeq 1/m_\rho$ 

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#### What about the 5d cutoff?

# From standard NDA considerations: $\Lambda_5 \simeq (16\pi^2) M_5 \left(\frac{4}{N_c}\right)^{2/3} = m_\rho \mathcal{O}(N_c^{1/3})$

The CS lowers the cutoff, but still well defined expansion

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We have 
$$\rho_S \sim L \sim 1/m_{\rho}$$
, expansion parameter:  
 $\frac{1}{\rho_S \Lambda_5} \rightarrow 0 \quad \text{for} \quad N_c \rightarrow \infty$ 

5d models provide Calculable Implementation of Skyrme's idea!!

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## Static Properties of the Nucleons:

	Experiment	AdS <sub>5</sub>	Deviation
M <sub>N</sub>	940 MeV	1130 MeV	+20%
$\mu_{S}$	0.44	0.34	-30%
$\mu_V$	2.35	1.79	-31%
ØА	1.25	0.70	-79%
$\sqrt{\langle r_{E,S}^2 \rangle}$	0.79 fm	0.88 fm	+11%
$\sqrt{\langle r_{E,V}^2 \rangle}$	0.93 fm	$\infty$	
$\sqrt{\langle r_{M,S}^2 \rangle}$	0.82 fm	0.92 fm	+12%
$\sqrt{\langle r_{M,V}^2 \rangle}$	0.87 fm	$\infty$	
$\sqrt{\langle r_A^2 \rangle}$	0.68 fm	0.76 fm	+12%
$\mu_p/\mu_n$	-1.461	-1.459	+0.1%

Significantly better agreement than original Skyrme model.

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What changes with pion masses? in progress with O. Domenech, G. Panico

Holography forces to introduce 5d scalar

	Experiment (MEV)	$AdS_5$ (MEV)	Deviation
$m_{\pi}$	$135 \mathrm{MeV}$	$134 \mathrm{MeV}$	0.6%
$m_{\pi(1300)}$	$1300 { m MeV}$	$1230 { m MeV}$	5.6%
$m_{ ho}$	$775 \mathrm{MeV}$	$783 { m MeV}$	1.0%
$m_\omega$	$782 \mathrm{MeV}$	$783 { m MeV}$	0.1%
$m_{a_1(1260)}$	$1230 { m MeV}$	$1320 { m MeV}$	7.6%
$m_{a_0(980)}$	$980 { m MeV}$	$1040 { m MeV}$	6.5%
$m_{f_0(980)}$	$980 { m MeV}$	$1040 { m MeV}$	6.5%
$f_{\pi}$	$92 { m MeV}$	$89 \mathrm{MeV}$	3.6%
$f_{ ho}$	$153 \mathrm{MeV}$	$149 \mathrm{MeV}$	2.7%
$\dot{f_\omega}$	$140 \mathrm{MeV}$	$149 \mathrm{MeV}$	6.4%
$g_{\rho\pi\pi}$	6.0	4.89	22.7%
$g_{\omega\pi\gamma}$	0.72	0.71	1.1%
$g_{\rho\pi\gamma}$	0.22	0.24	7.9%
$g_{\omega ho\pi}$	15.0	15.6	3.7%
RMSE			7.7%

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### Static Properties of the Nucleons (very preliminary)

	Experiment	$AdS_5$	Deviation
M <sub>N</sub>	940 MeV	$\sim 1070~{ m MeV}$	$\sim 14\%$
$\mu_{S}$	0.44	0.38	16%
$\mu_V$	2.35	$\sim 1.2$	$\sim 100\%$
ØА	1.25	$\sim$ 0.6	$\sim 100\%$
$\sqrt{\langle r_{E,S}^2 \rangle}$	0.79 fm	0.82 fm	4%
$\sqrt{\langle r_{E,V}^2 \rangle}$	0.93 fm	0.97 fm	4%
$\sqrt{\langle r_{M,S}^2 \rangle}$	0.82 fm	0.84 fm	3%
$\sqrt{\langle r_{M,V}^2 \rangle}$	0.87 fm	0.87 fm	0.5%
$\sqrt{\langle r_A^2 \rangle}$	0.68 fm	$\sim 0.6~\text{fm}$	$\sim 13\%$

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- ► 5d models mimic several expected features of large-N<sub>c</sub> QCD
- Provide a valid and "economic" description of QCD hadrons

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- Provide a valid and "economic" description of QCD hadrons
- Motivated by Bottom–Up phenomenological considerations: A relation with AdS/CFT? (seems unlikely)

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- 5d models mimic several expected features of large-N<sub>c</sub> QCD
- Provide a valid and "economic" description of QCD hadrons
- Motivated by Bottom–Up phenomenological considerations: A relation with AdS/CFT? (seems unlikely)
- Calculable Skryme model is automatically implemented

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## In Progress

The 5d model should be improved in several aspects:

- Add quark masses to the 5d skyrmion (almost done with O. Domenech and Giuliano Panico)
- 2. Include  $U(1)_A$  anomaly and  $\eta'$  mass (few literature, some ideas with G.Panico)

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#### The Sakai–Sugimoto Model

Equivalent to ours in the meson sector.

Only differs by the shape of the warp factor.

Perturbative Field Theory regime for  $N_c$ ,  $\lambda \to \infty$   $(M_{St} \sim \sqrt{\lambda}/L)$ .

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#### The Sakai–Sugimoto Model

Equivalent to ours in the meson sector.

Only differs by the shape of the warp factor.

Perturbative Field Theory regime for  $N_c, \lambda \to \infty$   $(M_{St} \sim \sqrt{\lambda}/L)$ .

The CS is subleading in  $1/\lambda$ :  $\gamma \sim 1/(M_{St}L)^2 \sim 1/\lambda$ .

For  $\gamma \ll 1$  the 5d skyrmion becomes small instanton:

$$E(
ho) \sim 
ho^2/L^2 + \gamma^2 L^2/
ho^2 \quad \Rightarrow \quad 
ho \sim \sqrt{\gamma} L \sim 1/M_{St}$$

Strong indication of non-calculability.

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