

# Asymptotic safety of gravity and the Higgs boson mass

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Based on: C. Wetterich, M. S., Phys. Lett. B683 (2010) 196

- Asymptotic safety versus renormalizability
- Asymptotically safe gravity
- Standard Model, observations, and asymptotic safety
- Higgs boson mass
- Conclusions

# Asymptotic safety versus renormalizability

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## Generic quantum field theory

- Take some field theory and write the most general Lagrangian.
- Compute all amplitudes in all orders of perturbation theory.
- Require that the theory is unitary, Lorentz - invariant, causal, etc - infinite number of conditions for infinite number of processes.
- Solve these consistency equations. Hopefully, the theory will be characterised by a finite number of essential parameters - coupling constants, making the predictions possible.

# RG approach

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Introduce dimensionless coupling constants  $g_i$  constants for all terms in the action:

$$g_i = \mu^D G_i, \quad G_i \text{ are dimensionfull in general}$$

$D$  is the dimension of coupling constant.

RG equations: from requirement that physical amplitudes are  $\mu$ -independent,

$$\mu \frac{\partial g_i}{\partial \mu} = \beta_i(g)$$

# Different possibilities

- Renormalizable asymptotically free theories – Gaussian UV fixed point: essential couplings  $g_i \rightarrow 0$  at  $\mu \rightarrow \infty$ . The number of these couplings is finite - only operators with dimension  $\leq 4$  are allowed.
- Asymptotically safe theories – non-Gaussian UV fixed point  $g^* \neq 0$ :  $\beta_i(g^*) = 0$ . If the dimensionality of the critical surface in the space of coupling constants (which points are attracted to  $g^*$ ) is finite, the theory is predictable.

# Known solutions

- Asymptotically free theories:
  - QCD
  - Certain GUTS
  - Renormalizable theories in 2d and 3d
- Asymptotically safe, but non-renormalizable theories
  - Scalar field theory in 3d at Wilson-Fischer fixed point (critical surface is 2-dimensional)
  - Non-linear  $\sigma$  model in 3d
  - Complete theory of pions and nucleons in 4d

The Standard Model is neither asymptotically free nor asymptotically safe!

# Methods to study asymptotic safety

- $\epsilon$  - expansion
- Lattice simulations
- Functional renormalisation group:

$$\mu \partial_\mu S_\mu = \frac{1}{2} \text{STr} \left[ \left( S_\mu^{(2)} + \mathcal{R}_\mu \right)^{-1} \mu \partial_\mu \mathcal{R}_\mu \right] .$$

where

$S_\mu \rightarrow S_{class}$  for  $\mu \rightarrow \infty$ ,

$S_\mu \rightarrow S_{eff}$  for  $\mu \rightarrow 0$ ,

$\mathcal{R}_\mu$  is the “window” function.

Conjecture **Weinberg '79**: Gravity may be asymptotically safe.

$\epsilon$ -expansion argument -

$$S_G = -\frac{1}{16\pi G_0} \int d^D x \sqrt{g} R$$

$$G(\mu) = \mu^{D-2} G_0, \quad \mu \frac{d}{d\mu} G(\mu) = (D-2)G(\mu) - bG^2(\mu) .$$

Fixed point:

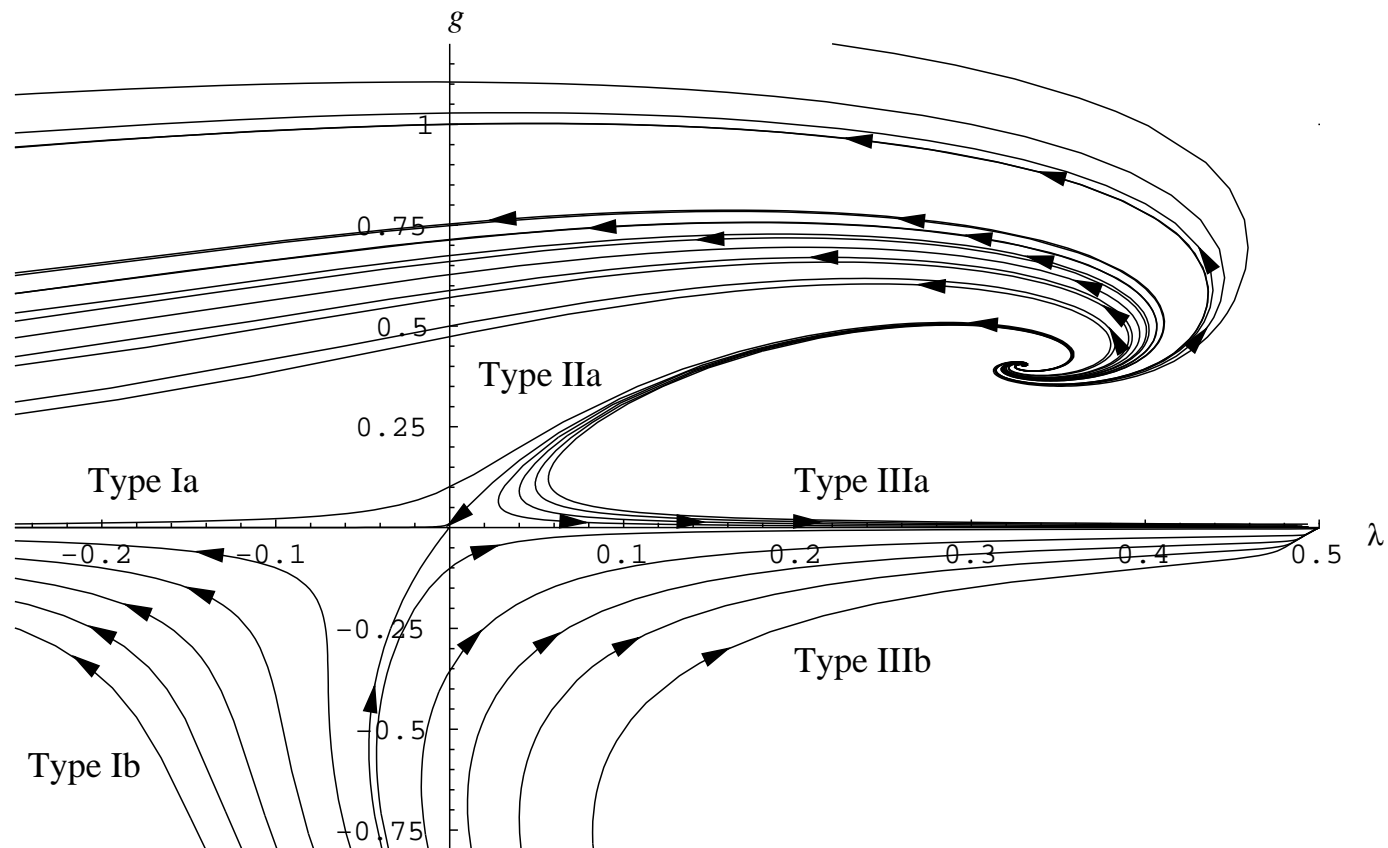
$$G^* = \frac{D-2}{b}, \quad G_0^*(\mu) = \frac{G^*}{\mu^{D-2}} \rightarrow 0 \text{ if } \mu \rightarrow \infty$$

Computations give  $b > 0$ . **Gastmans et al '77, Christensen and Duff '77, Kawai and Ninomiya '90, Percacci '06,...**



Functional RG analysis - Reuter '96, Percacci et al, Niedermaier '09, ...

$$S_{\text{EH}} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \{-R + 2\Lambda\} ,$$



Reuter and Saueressig '02

# Extra evidence

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- Higher derivative gravity – Stelle '77, Fradkin & Tseytlin '82, Avramidi & Barvinsky '85,...
- Large N (matter fields) expansion – Tomboulis '77, '80, Smolin '82, Percacci '06,...
- Perturbation theory – Niedermaier, '09

What if indeed gravity is asymptotically safe?

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Any predictions for particle physics and LHC?

# Possible consequence: Electroweak theory + Gravity is a final theory

Experimental evidence for physics beyond the SM

- i. Neutrino masses and oscillations
- ii. Dark matter
- iii. Baryon asymmetry of the Universe
- iv. Inflation

require only a modest extension of the SM ( $\nu$ MSM) by 3 singlet right-handed fermions (needed for **i-iii**) with masses in keV - GeV area, and non-minimal coupling of the Standard Model Higgs field to Ricci scalar (needed for **iv**).

# Realisation: $\nu$ MSM

SM fermions

quarks

left	u	d	c	s	t	b
right	u	d	c	s	t	b
left	$\nu_e$	e	$\nu_\mu$	$\mu$	$\nu_\tau$	$\tau$
right		e		$\mu$		$\tau$

leptons

$\nu$ MSM fermions

quarks

left	u	d	c	s	t	b
right	u	d	c	s	t	b
left	$\nu_e$	e	$\nu_\mu$	$\mu$	$\nu_\tau$	$\tau$
right	$N_e$	e	$N_\mu$	$\mu$	$N_\tau$	$\tau$

leptons

**Role** of  $N_e$  with mass in keV region: dark matter

**Role** of  $N_\mu$ ,  $N_\tau$  with mass in 100 MeV – GeV region: “give” masses to neutrinos and produce baryon asymmetry of the Universe

**Role** of the Higgs: give masses to quarks, leptons,  $Z$  and  $W$  and inflate the Universe.

# To be true: all the couplings of the SM must be asymptotically safe or asymptotically free

Problem for:

- U(1) gauge coupling  $g_1$ ,  $\mu \frac{dg_1}{d\mu} = \beta_1^{\text{SM}} = \frac{41}{96\pi^2} g_1^3$

- Scalar self-coupling  $\lambda$ ,  $\mu \frac{d\lambda}{d\mu} = \beta_\lambda^{\text{SM}} =$

$$= \frac{1}{16\pi^2} \left[ (24\lambda + 12h^2 - 9(g_2^2 + \frac{1}{3}g_1^2))\lambda - 6h^4 + \frac{9}{8}g_2^4 + \frac{3}{8}g_1^4 + \frac{3}{4}g_2^2g_1^2 \right]$$

- Fermion Yukawa couplings, t-quark in particular  $h$ ,  $\mu \frac{dh}{d\mu} = \beta_h^{\text{SM}} =$

$$= \frac{h}{16\pi^2} \left[ \frac{9}{2}h^2 - 8g_3^2 - \frac{9}{4}g_2^2 - \frac{17}{12}g_1^2 \right]$$

Landau pole behaviour

# Gravity contribution to RG running

Let  $x_j$  is a SM coupling. Gravity contribution to RG:

$$\mu \frac{dx_j}{d\mu} = \beta_j^{\text{SM}} + \beta_j^{\text{grav}} .$$

On dimensional grounds

$$\beta_j^{\text{grav}} = \frac{a_j}{8\pi} \frac{\mu^2}{M_P^2(\mu)} x_j .$$

where

$$M_P^2(\mu) = M_P^2 + 2\xi_0\mu^2 ,$$

with  $M_P = (8\pi G_N)^{-1/2} = 2.4 \times 10^{18}$  GeV,  $\xi_0 \approx 0.024$

from a numerical solution of FRGE



- The couplings are **not** in  $\overline{MS}$  scheme
- The couplings are **not** in MOM scheme
- Pretty vague definition based on physical scattering amplitudes at large momentum transfer - never actually worked out in details

Thus, computations of  $a_j$  are ambiguous and controversial.

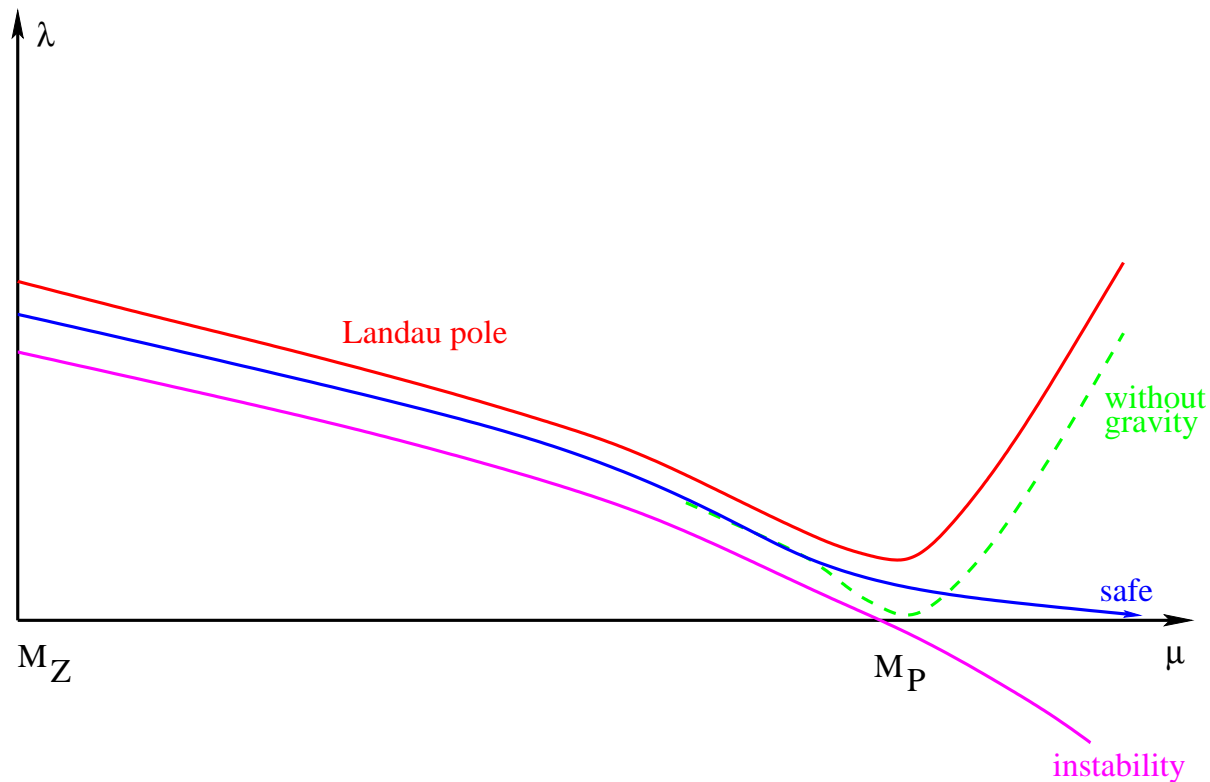
Still, even without exact knowledge of  $a_j$  a lot can be said about the Higgs mass

Robinson and Wilczek '05, Pietrykowski '06, Toms '07&'08, Ebert, Plefka and Rodigast '07, Narain and Percacci '09, Daum, Harst and Reuter '09, Zanusso et al '09, ...

- Most works get for gauge couplings a universal value  
 $a_1 = a_2 = a_3 < 0$ : U(1) gauge coupling get asymptotically free in asymptotically safe gravity
- $a_\lambda \simeq 2.6 > 0$  according to Percacci and Narain '03 for scalar theory coupled to gravity
- $a_h > < 0$  ?? The case  $a_h > 0$  is not phenomenologically acceptable - only massless fermions are admitted

Suppose that indeed  $a_1 < 0$ ,  $a_h < 0$ ,  $a_\lambda > 0$ . Then the Higgs mass is predicted with theoretical uncertainty  $\simeq \pm 2.2$  GeV

$$m_H = \left[ 126.3 + \frac{m_t - 171.2}{2.1} \times 4.1 - \frac{\alpha_s - 0.1176}{0.002} \times 1.5 \right] \text{ GeV} ,$$



Possible understanding of the amazing fact that  $\lambda(M_P) = 0$  and  $\beta_\lambda^{\text{SM}}(M_P) = 0$  simultaneously at the Planck scale.

To decrease uncertainty: (the LHC accuracy can be as small as **200 MeV!**)

- Measure better t-quark mass (present error in  $m_H$  due to this uncertainty is  $\simeq 4$  GeV)
- Measure better  $\alpha_s$  (present error in  $m_H$  due to this uncertainty is  $\simeq 1.5$  GeV)
- Compute two-loop EW corrections to pole -  $\overline{MS}$  matching for the Higgs mass (has never been done)
- Compute 3-loop running of all couplings of the Standard Model (has never been done)

If done, the uncertainty will be reduced to  $\sim 0.5$  GeV, due to unremovable non-perturbative contribution  $\sim \Lambda_{QCD}$  to top quark mass.

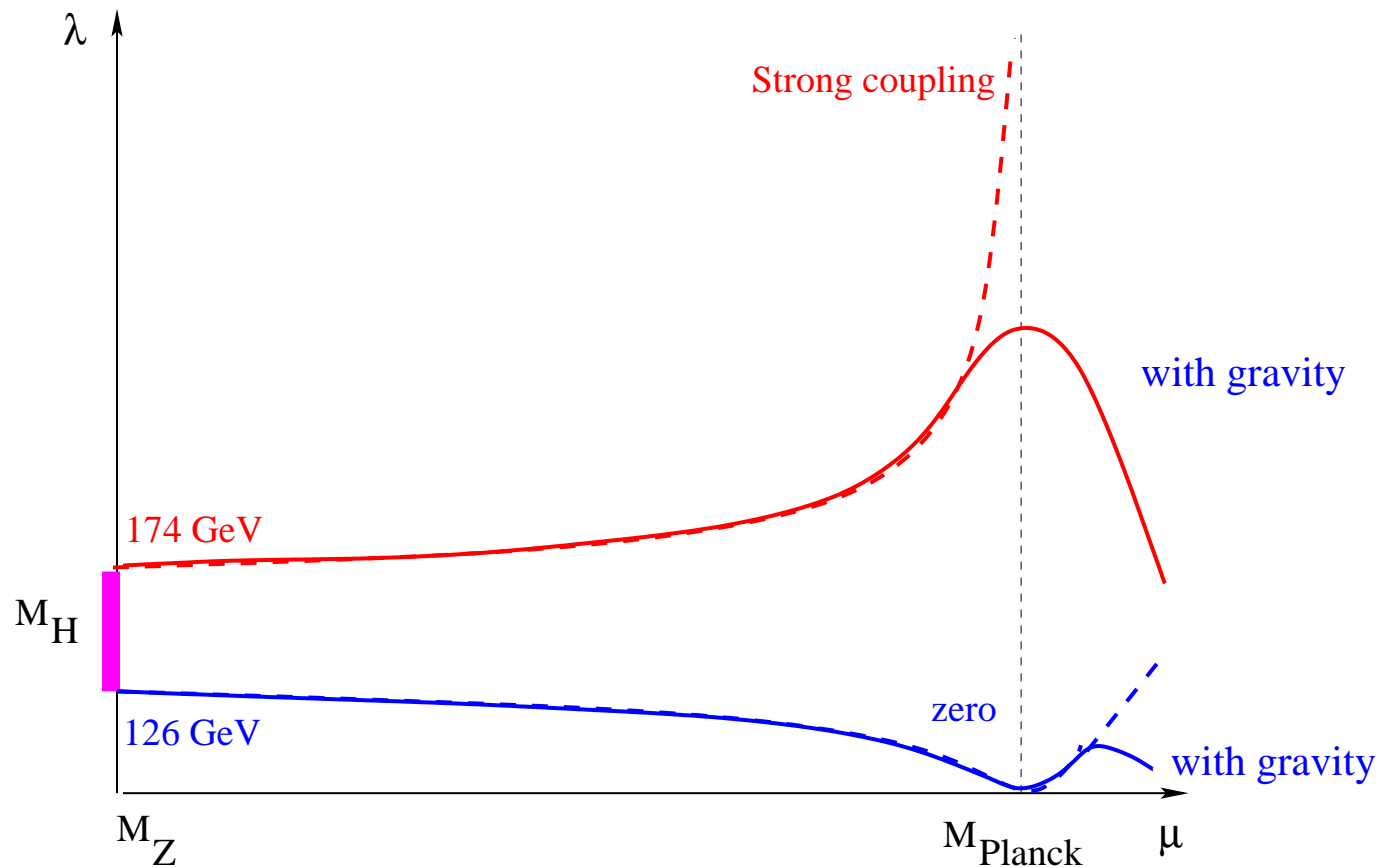
Prediction is quite model independent and may be valid for GUTs, higher dimensional theories, etc.

It stays approximately true provided:

- We have SM running of couplings up to the high energy scale
- $a_\lambda$  is relatively large and positive
- If there are other light particles, the number for  $m_H$  will be different.

Suppose that  $a_1 < 0$ ,  $a_h < 0$ ,  $a_\lambda < 0$ . Then the Higgs mass is predicted with theoretical uncertainty  $\simeq 50$  GeV

$$126 \text{ GeV} < m_H < 174 \text{ GeV}$$



# Asymptotic safety and the Fermi scale

Assumed: the Fermi scale is fixed to its experimental value. Gravity contribution to running of the mass parameter in the Higgs-potential  $m^2(\mu)$ :

$$\mu \frac{\partial}{\partial \mu} m^2 = A_m m^2 ,$$

If:  $A_m > 2$  the dimensionless ratio  $m^2/\mu^2$  is attracted to zero, leading, possibly, to understanding why  $G_F \gg G_N$ .

Percacci '03, Narain '09:  $A_m = 1.83 < 2$  for scalar-Gravity system.

SM+Gravity - not known

# Conclusions

If gravity is asymptotically safe then the possible outcome of the LHC experiments is:

- Higgs and nothing else
- $m_H \simeq 126 \text{ GeV}$  (for central values of  $m_t$  and  $\alpha_s$ ) if, as some computations show,  $a_\lambda > 0$
- $126 \text{ GeV} < m_H < 174 \text{ GeV}$  if  $a_\lambda < 0$
- Waiting time  $\sim 6$  years (?)
- Asymptotic safety may shed light to the smallness of the Fermi scale