

Observing GW signature from supermassive BH mergings

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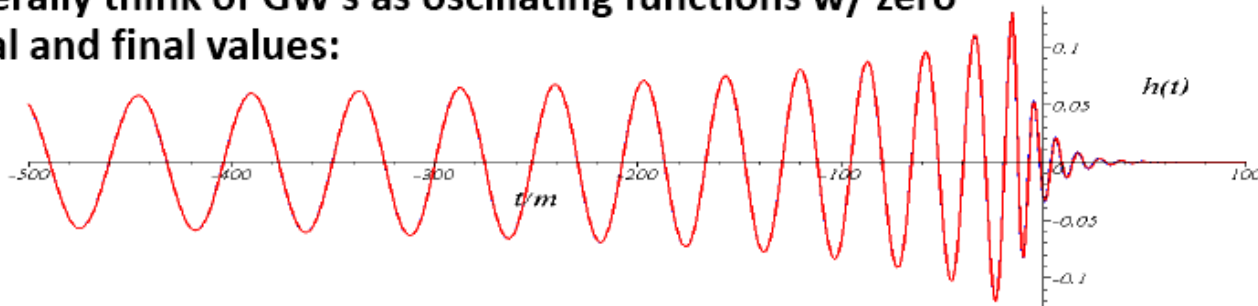
0909.0742, MNRAS 2010

Outline

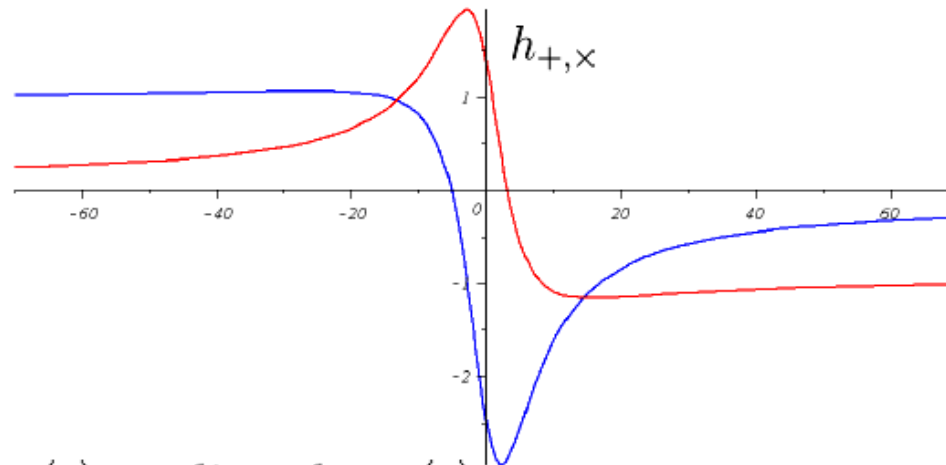
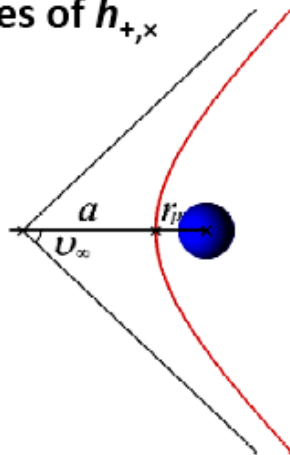
- GW bursts with memory
- SMBH mergings
- Signatures in pulsar timing
- S/N ratio
- Expected event rate
- Conclusions

GW bursts with memory

- Generally think of GW's as oscillating functions w/ zero initial and final values:

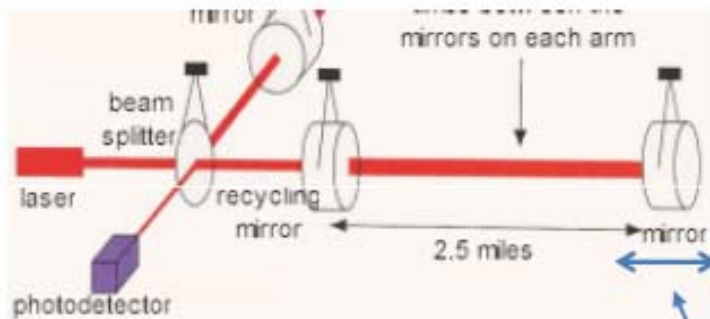


- But some sources exhibit differences in the initial & final values of $h_{+, \times}$



$$\Delta h_{+, \times}^{\text{mem}} = \lim_{t \rightarrow +\infty} h_{+, \times}(t) - \lim_{t \rightarrow -\infty} h_{+, \times}(t)$$

- Ideal GW detector (‘free masses’) would have permanent displacement (‘memory’)

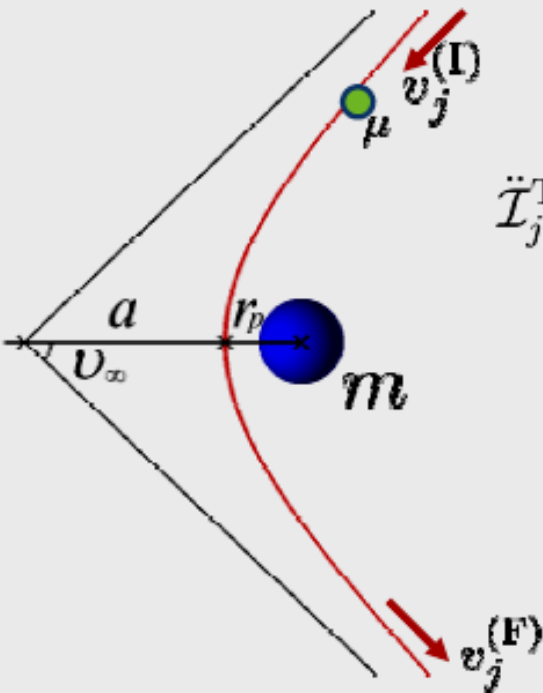


$$\delta x(t) = \frac{L}{2} h_+(t) \Rightarrow \delta x^{\text{mem}} = \frac{L}{2} h_+^{\text{mem}}$$

- Build-up of the displacement is measurable (difficult by ground-based LIGO, but can be done from space by LISA)

GWM: linear effect

- Non-oscillatory change of quadrupole and higher multipole moments (Zeldovich & Polnarev 74, Braginsky & Grishchuk 78, Braginsky & Thorne 87). For example, gravitational scattering (hyperbolic orbit)



$$h_{jk}^{\text{TT}} \approx \frac{2}{R} \ddot{\mathcal{I}}_{jk}^{\text{TT}} \quad \mathcal{I}_{jk}^{\text{TT}} = \mu [x_j x_k]^{\text{TT}}$$

$$\ddot{\mathcal{I}}_{jk}^{\text{TT}} = \mu [x_j \ddot{x}_k + \ddot{x}_j x_k + 2\dot{x}_j \dot{x}_k]^{\text{TT}}$$

$$= 2\mu \left[\dot{x}_j \dot{x}_k - \frac{M}{r^3} x_j x_k \right]^{\text{TT}}$$

$$\ddot{x}_j = -\frac{M}{r^3} x_j$$

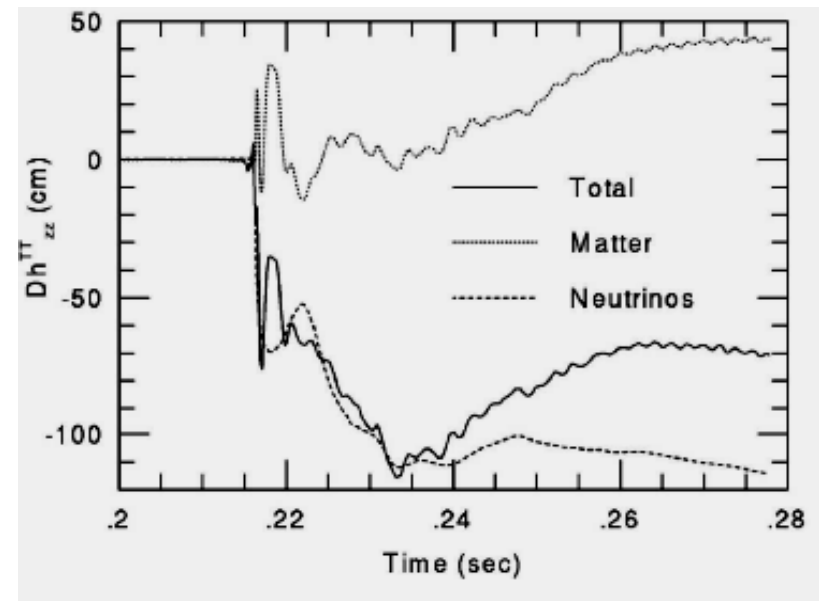
Turner 77,
Velocities before and after scattering are the same but sign is opposite

$$\Delta h_{jk}^{\text{TT}} = \frac{4\mu}{R} \Delta [\dot{x}_j \dot{x}_k]^{\text{TT}}$$

- For any system of N gravitationally unbound bodies with velocities v_A before and after GW burst emission (Braginsky & Thorne 87, Thorne 92)

$$\Delta h_{jk}^{\text{TT}} = \Delta \sum_{A=1}^N \frac{4M_A}{R\sqrt{1-v_A^2}} \left[\frac{v_A^j v_A^k}{1-v_A \cdot N} \right]^{\text{TT}}$$

- Examples: hyperbolic orbits (Turner 77), asymmetric neutrino emission (Epstein 78), asymmetric SN explosions (Burrows & Hayes 96, Ott 08), GRB jets...



Nonlinear effect (Cristodoulou memory)

Cristodoulou 91, Blanchet & Damour 92

- Contribution to the distant GW field sourced by the emission of GWs
- Recall previous form of the Einstein's equations:

$$\square h_{\alpha\beta} = 16\pi \det(g_{\mu\nu}) T_{\alpha\beta} + \mathcal{F}[h, h]$$

Grav'l wave stress-energy tensor...

$$T_{\alpha\beta}^{\text{gw}} \propto \frac{dE^{\text{gw}}}{dt d\Omega} \sim O(h^2)$$

$$\ddot{\mathcal{I}}_{jk} \rightarrow \ddot{\mathcal{I}}_{jk} + U_{jk}^{\text{gw}}$$

...contributes to the changing multipole moments...

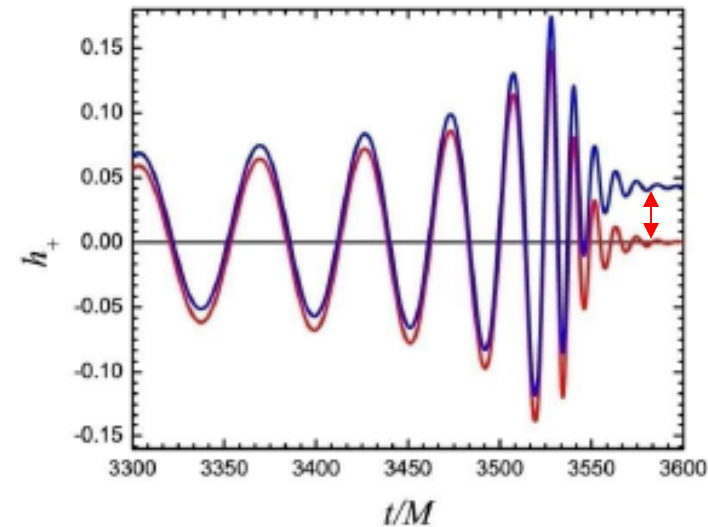
(Favata 09)

$$h_{jk}^{\text{TT}} \approx \frac{2}{R} \ddot{\mathcal{I}}_{jk}^{\text{TT}}$$

...which determines the GW field...

$$\Delta h^{(\text{mem})} \sim \frac{\Delta E^{\text{gw}}}{R}$$

...which has a slowly-growing, non-oscillatory piece related to the radiated GW energy.



- Hereditary nature: memory piece of the GW field depends on the entire history:

$$\delta h_{jk}^{\text{TT}} = \frac{4}{R} \int_{-\infty}^{T_R} dt' \left[\int \frac{dE^{\text{SW}}}{dt' d\Omega'} \frac{n'_j n'_k}{(1 - \mathbf{n}' \cdot \mathbf{N})} d\Omega' \right]^{\text{TT}} \quad (T_R \text{ is retarded time})$$

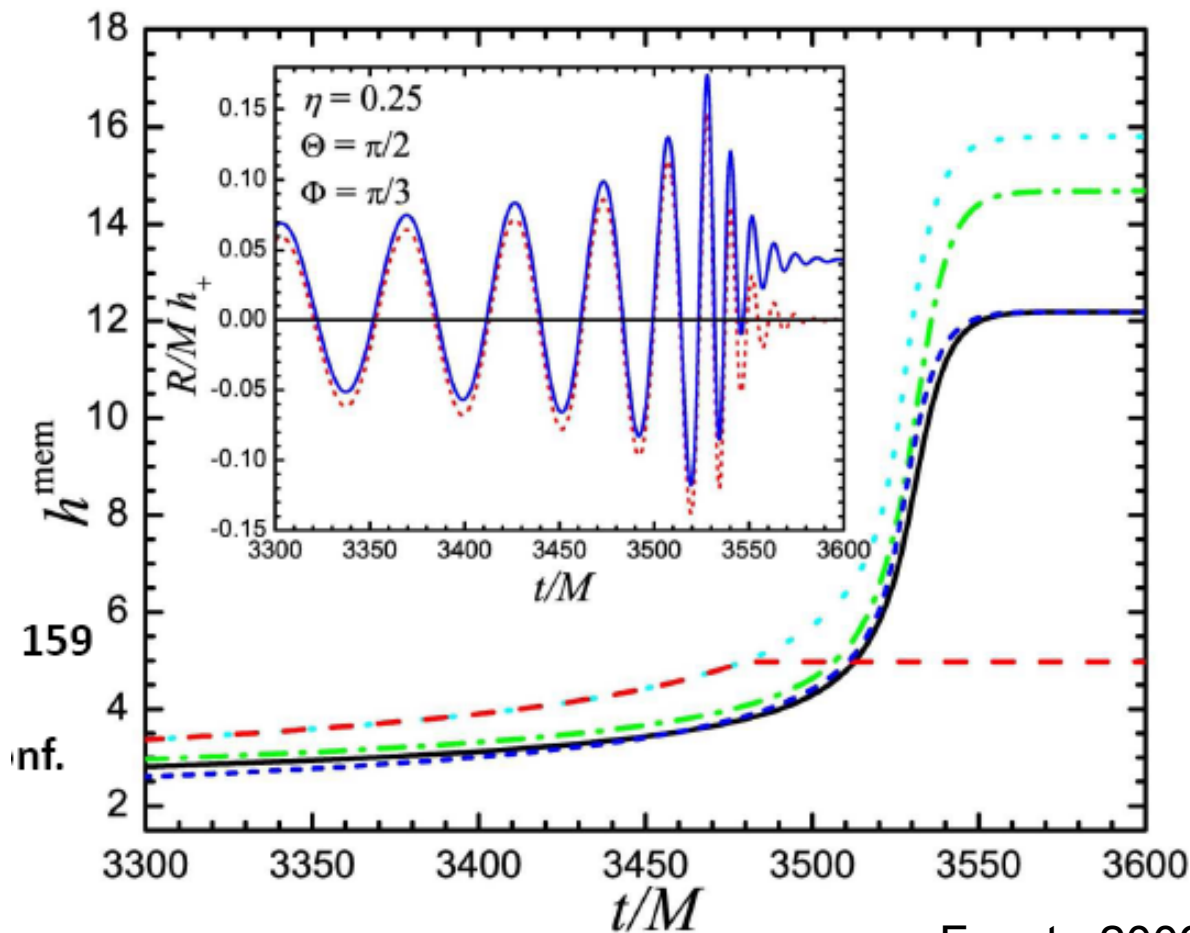
- Similar to linear memory can be interpreted as arising from changes in the mass quadrupole moment of the system during emission of individual gravitons

(Thorne 92) with energies $E_A = M_A / (1 - v_A^2)^{1/2}$

and velocities $v_a^j = c n_a^j$

GWM in binary BH mergers

- Quasi-circular orbits: $h_x=0$, $h_+ \neq 0$



Depends on the entire history, but roughly

$$h^{\text{mem}} \sim h^{\text{insp}}$$

$$h_+^{mem} = \frac{\eta M h}{384\pi R} \sin^2 \theta (17 + \cos^2 \theta), \quad M = M_1 + M_2, \quad \eta = \frac{M_1 M_2}{M^2}$$

θ - angle between orbital ang. momentum and line of sight

$$h = \frac{16\pi}{\eta} \left(\frac{\Delta E_{GW}}{M} \right)$$

$$\langle h_+^{mem} \rangle = \frac{69}{8} \left(\frac{\Delta E_{GW}}{24R} \right) \approx \frac{\Delta E_{GW}}{3R}$$

From numerical simulations (Reisswig et al 09):

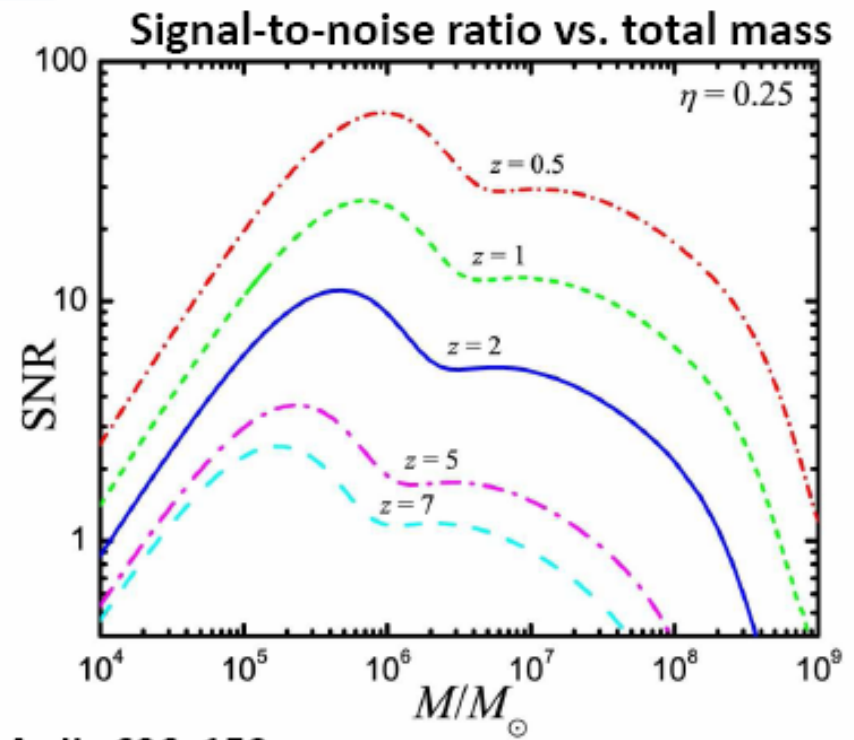
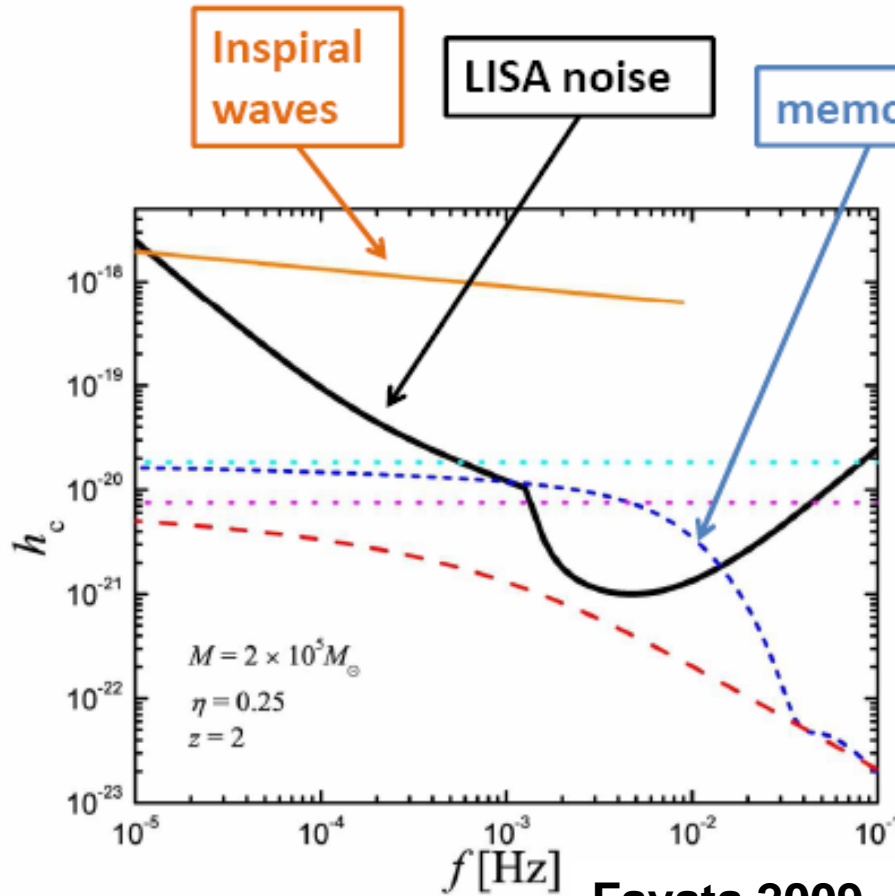
$$\Delta E_{GW} \approx (3.6 - 10\%) M \quad \Rightarrow$$

$$h^{mem} \approx 5 \times 10^{-16} \left(\frac{m}{10^8 M_\odot} \right) \left(\frac{1 \text{ Gpc}}{R} \right)$$

Detectability of the memory:

SMBH mergers

- will be difficult to observe w/ Advanced LIGO
- likely to be visible by LISA out to redshift $z \lesssim 2$



Favata 2009, ApJL, 696, 159

Detection of GWM by pulsar timing

- GWM leaves unique signature in pulsar timing: linear growth of rms residuals with time

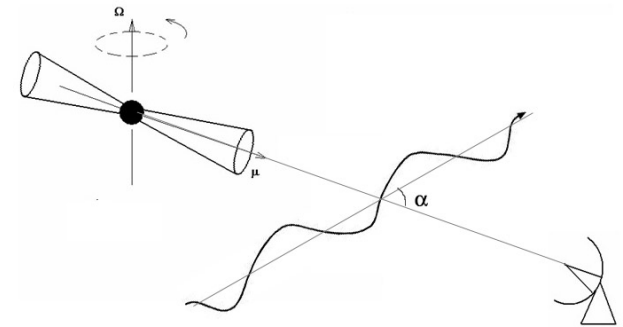
$$S_{mem} \sim h_+^{mem} T_{obs}$$

$$S_{insp} \sim h^{insp} / \omega_{insp}$$

$$\omega_{insp}^{-1} \approx 7.5 \times 10^3 (M / 10^8 M_\odot) s$$

$$T_{obs} \approx 10 \text{ yrs}$$

$$\frac{SNR_{GWM}}{SNR_{insp}} \sim \frac{h_+^{mem} T_{obs} \omega_{insp}}{h^{insp}} \sim T_{obs} \omega_{insp} \sim 2 \times 10^4$$



Open prospects for detection!

Pulsar frequency modulation (Sazhin 78, Detweiler 79)

$$\frac{\Delta \nu}{\nu_0} = \frac{1}{2} \int_0^D d\lambda \left(e^i e^j \frac{\partial h^{ij}}{\partial t} \right) \Big|_{path}$$

Timing residuals

$$s(t) = \int_0^t d\tau \frac{\Delta \nu(\tau)}{\nu_0}$$

For plain gravitational wave

$$h_{ij}(x^i, t) = h_+(t - n_i x^i) p_{ij}^+ + h_\times(t - n_i x^i) p_{ij}^\times$$

$$\frac{\Delta \nu(t)}{\nu_0} = \frac{1}{2} (1 + \mu) \left\{ \begin{aligned} & [h_+(t) \cos 2\phi + h_\times(t) \sin 2\phi] - \\ & [h_+(t - D(1 - \mu)) \cos 2\phi + h_\times(t - D(1 - \mu)) \sin 2\phi] \end{aligned} \right\}$$

Physical sense: frequency variation is determined by difference between the GW strength at the place and time of observations (first [.]) and its strength at the site and time of signal emission

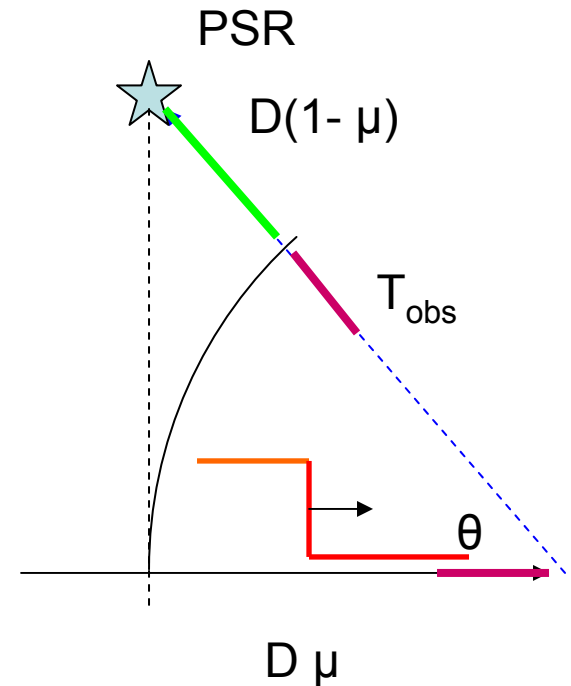
For GWBM from SMBH $h_{\times} = 0$,
 $h_{+}(t - D(1 - \mu)) = 0$ at the site of PSR
 if $D(1 - \mu) > T_{obs}$.

Typically $D \sim \text{kpc} \gg T_{obs} \sim 10 \text{ yrs}$

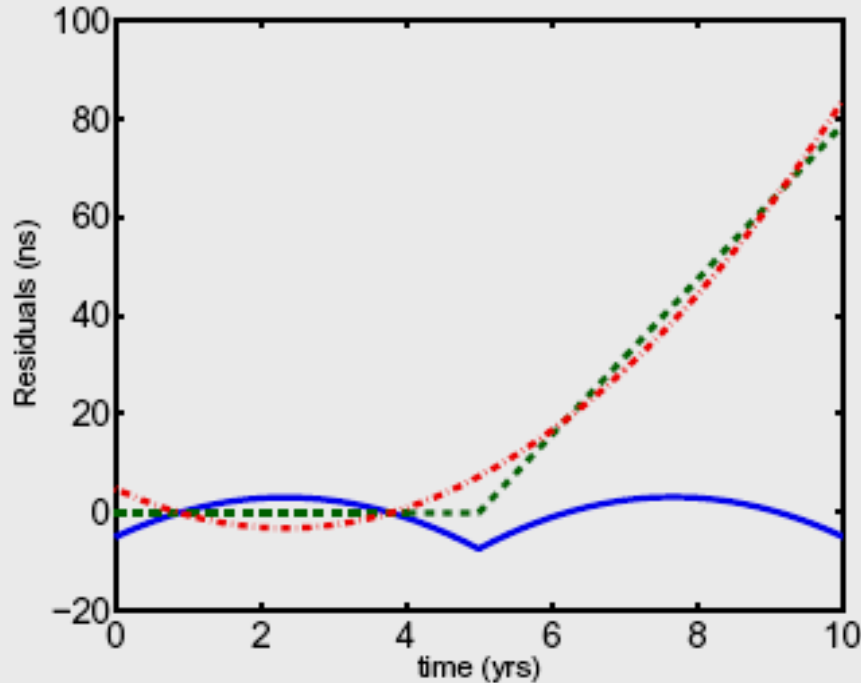
Net result:

$$s(t) = \frac{1}{2} (1 + \mu) \int_0^t d\tau h_{+}(\tau) \cos 2\phi$$

$$\mu = \cos \theta$$



Timing residuals



- Calculate expected signal
- Extract quadratic fit to obtain postfit residuals
- Check whether they can be measured at a given SNR by pulsar timing array (PTA)

Pshirkov, Baskaran, PK, 2010 MNRAS

Detectability: 1. SNR

Residuals (t) = signal residuals (t) + noise (t)

Gaussian stationary noise uncorrelated for each pulsar

$$\overline{n_{\alpha}(t_i)n_{\beta}(t_j)} = \sigma_n^2 \delta_{ij} \delta_{\alpha\beta}$$

PTA includes $N_{\alpha} = 20$ PSRs, $T_{\text{obs}} \approx 10$ yrs,

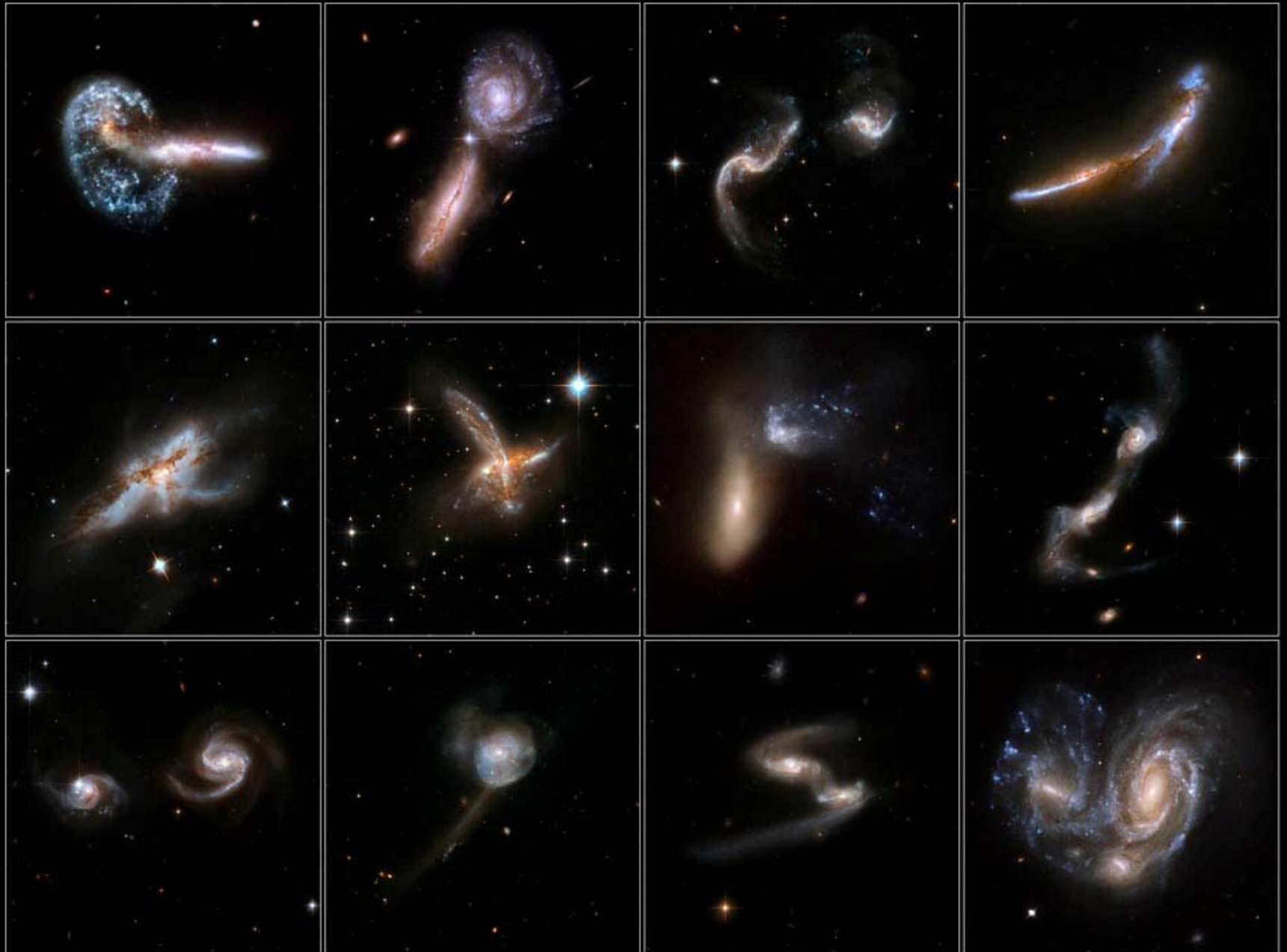
of individual observations $N_t = 250$

Expected sensitivity: $\sigma_n = 100$ ns

$$\text{SNR}_{\text{GWM}} \approx 1.6 \left(\frac{h^{\text{mem}}}{10^{-15}} \right) \left(\frac{N_t}{250} \right)^{1/2} \left(\frac{N_t}{250} \right)^{1/2} \left(\frac{T_{\text{obs}}}{10 \text{ yrs}} \right) \left(\frac{100 \text{ ns}}{\sigma_n} \right)$$

Interacting Galaxies

Hubble Space Telescope • ACS/WFC • WFPC2



Detectability: 2. Rate of events

- Rough estimate: $\sim 10^{-3} \text{ Gpc}^{-3} \text{ yr}^{-1}$ as major galaxy mergers (Concelice et al 09) $\rightarrow \sim 0.5$ event per year up to $z \sim 0.5$. Consistent with numerical simulations
- More accurate analysis: from observed SMBH mass function $N(>M) = 0.07(M/10^7)^{-2}$ and galaxy merger rate evolution $n(z) \sim (1+z)^{(2 \dots 3)}$

$$N \approx 0.1 \left(\frac{N_t}{250} \right) \left(\frac{N_\alpha}{20} \right) \left(\frac{T_{obs}}{10 \text{ yrs}} \right)^3 \left(\frac{100 \text{ ns}}{\sigma_n} \right)^2 \left(\frac{3}{SNR} \right)$$

Conclusions

- GW bursts with memory leave unique imprint in pulsar timing residuals
- Gravitational bursts with memory and amplitude $h \sim (1.5-2) \times 10^{-15}$ can be potentially detected in 10yrs of Pulsar Timing Array observations at the SNR=3 at the current timing noise level 100 ns.
- With account for observed SMBH mass function and evolution of galaxy major mergers with redshift, PTA detection rate can be around 1 in a 10yrs run
- For merging SMBH, this method is complimentary to LISA for $M > 10^8 M_{\odot}$ to which LISA sensitivity decreases



THANK YOU!

10.06.2010

Quarks-2010