

Outline

- ▶ $SU(2)_{brane} \times U(1)_{bulk}$ model, Z boson localization.
- ▶ Processes with particle escape from brane. Cross-section computation.
- ▶ Comparison with ADD prediction.

$SU(2)_{brane} \times U(1)_{bulk}$

- ▶ (4+1+n) space time

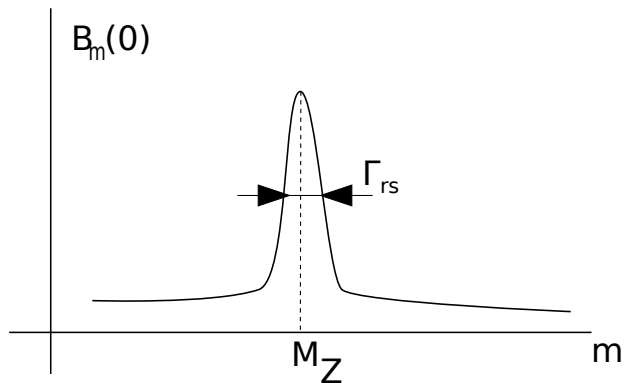
Action. $SU(2)$ and Higgs fields on the brane, $U(1)$ in bulk.

$$S = \int d^x dz \Pi_i \frac{d\theta_i}{2\pi R_i} \sqrt{|g|} \left(\left(-\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \right. \right. \\ \left. \left. + (D\phi)^\dagger D\phi - \frac{\lambda}{4} \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2 \right) \delta(z) \right. \\ \left. - \frac{1}{4} B_{\bar{\mu}\bar{\nu}} B_{\bar{\gamma}\bar{\delta}} g^{\bar{\mu}\bar{\gamma}} g^{\bar{\nu}\bar{\delta}} \right) \quad (1)$$

Effective action

$$S_{eff} = \int d^4 x dm \left(-\frac{1}{4} B_{\mu\nu} B_{\mu\nu} + \frac{m^2}{2} B_\mu B_\mu \right) + \\ + \int d^4 x dm B_m(0) \bar{\psi} \gamma^\mu \psi B_\mu \quad (2)$$

Z boson invisible width



Width of particle escape from the brane

$$\Gamma_{rs} = \frac{\pi(M_Z^2 - M_W^2)}{\Gamma(\frac{n}{2})\Gamma(\frac{n}{2} + 1)M_Z} \left(\frac{m_z}{k}\right)^n \quad (3)$$

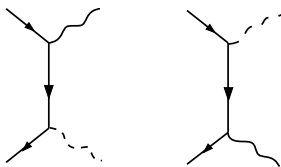
Limits for k

Limits for k from Z invisible width

$\Gamma_{invis} = 0.4990 \pm 0.0003$ GeV (PDG)

- ▶ $n = 1, k > 1.2 \cdot 10^4$ TeV
- ▶ $n = 2, k > 40$ TeV
- ▶ $n = 3, k > 5$ TeV
- ▶ $n = 4, k > 1.6$ TeV

Collider processes $e^+e^- \rightarrow \gamma + B_{bulk}$



In the rest frame of fermions the differential cross-section

$$\frac{d\sigma}{dx dy} = \left(1 + \frac{(\sin^2 \theta_w - 0.5)^2 + \sin^4 \theta_w}{2 \sin^2 \theta_w \cos^2 \theta_w} \right) \frac{4\pi\alpha^2}{n\Gamma^2(\frac{n}{2})} \cdot \frac{1}{s} \left(\frac{\sqrt{s}}{2k} \right)^n \cdot f(x, y) \quad (4)$$

$$f(x, y) = \frac{(1-x)^{\frac{n}{2}-1}}{x} \cdot \frac{(2-x)^2 + x^2 y^2}{1-y^2} \quad (5)$$

here $x = 2E/\sqrt{s}$, $y = \cos \theta$

The ADD prediction

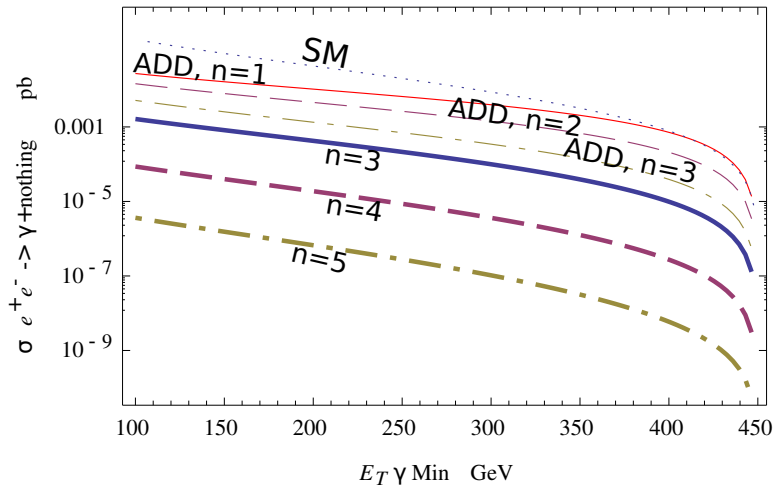
the ADD cross section to graviton and photon computed by G.F. Giudice, R. Rattazzi and J.D. Wells (hep-ph/9811291)

$$\frac{d^2\sigma}{dx_\gamma d\cos\theta}(e^+e^- \rightarrow \gamma G) = \frac{\alpha}{64} \cdot \frac{2\pi^{n/2}}{\Gamma(n/2)} \cdot \left(\frac{\sqrt{s}}{M_D}\right)^{n+2} \frac{1}{s} f(x_\gamma, \cos\theta) \quad (6)$$

$$f(x, y) = \frac{2(1-x)^{\frac{n}{2}-1}}{x(1-y^2)} \cdot [(2-x)^2(1-x+x^2) - 3y^2x^2(1-x) - y^4x^4]. \quad (7)$$

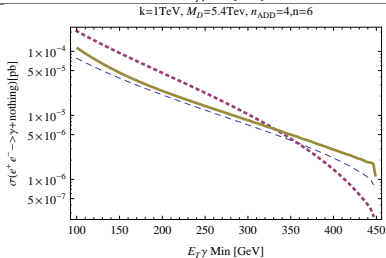
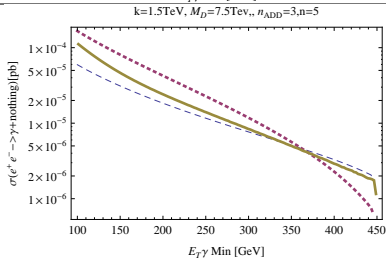
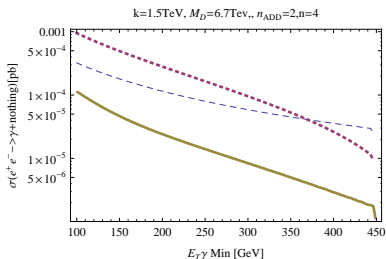
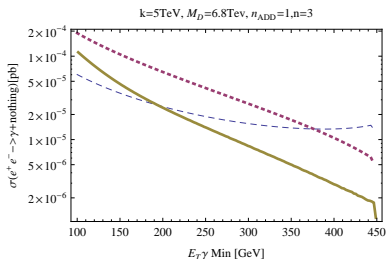
Comparison with ADD prediction

$k=4\text{TeV}, M_D=4\text{TeV}$



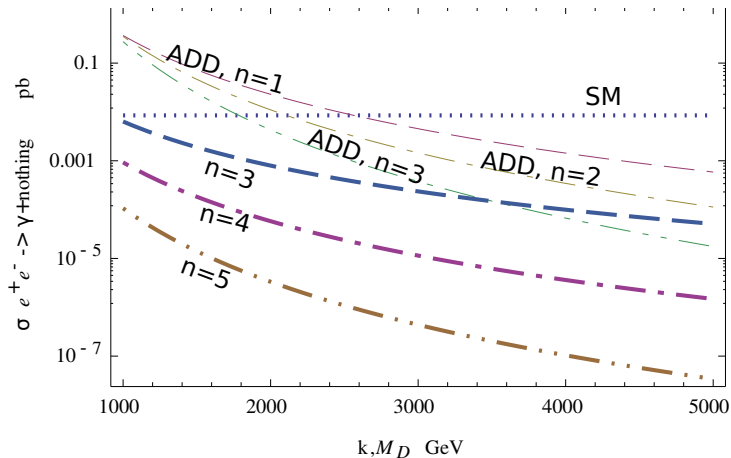
$\sqrt{s} = 1\text{TeV}, E_\gamma^T > 100\text{ GeV}, E_\gamma < 450\text{ GeV}$, integrated over angle between photon and fermion momentum.

Comparison with ADD prediction



$E_\gamma^T > 100 \text{ GeV}, E_\gamma < 450 \text{ GeV}$, integrated over angle between photon and fermion momentum $84^\circ \leq \theta \leq 90^\circ$

Comparison with ADD prediction



$\sqrt{s} = 1 \text{ TeV}$, $300 \text{ GeV} < E_\gamma < 450 \text{ GeV}$, integrated over angle between photon and fermion momentum

$E_\gamma^T < 300 \text{ GeV}$.

- ▶ It is possible to distinguish this model and ADD.
- ▶ $n = 3$, $k > 5\text{TeV}$ is interesting for future colliders.

Thank for your attention!

$SU(2)_{brane} \times U(1)_{bulk}$

Metric

$$ds^2 = e^{-2k|z|} \left(dt^2 - d\vec{x}^2 - \sum_i d\theta_i^2 \right) - dz^2 \quad (8)$$

Action

$$S = \int d^x dz \Pi_i \frac{d\theta_i}{2\pi R_i} \sqrt{|g|} \left(\left(-\frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \right. \right. \\ \left. \left. + (D\phi)^\dagger D\phi - \frac{\lambda}{4} \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2 \right) \delta(z) \right. \\ \left. - \frac{1}{4} B_{\bar{\mu}\bar{\nu}} B_{\bar{\gamma}\bar{\delta}} g^{\bar{\mu}\bar{\gamma}} g^{\bar{\nu}\bar{\delta}} \right) \quad (9)$$

Interaction with fermions

$$S_{int} = \int d^4 x dz \delta(z) \bar{\psi} \gamma^\mu \psi A_\mu \quad (10)$$

Classical equations of motion after symmetry breaking, unitary gauge

$$\begin{aligned} \frac{1}{a^2} \partial_\mu \partial_\mu B_\nu - \partial_z^2 B_\nu + (n+2)k \text{sign}(z) \partial_z B_\nu &= \\ &= \frac{2v^2 g'}{8} \sqrt{\frac{2}{kn}} \left(g' \sqrt{\frac{2}{kn}} B_\nu - g A_\nu^3 \right) \delta(z) \end{aligned} \quad (11)$$

$$\partial_\mu F_{\mu\nu} + \frac{2gv^2}{8} \left(g A_\nu^3 - g' \sqrt{\frac{2}{kn}} B_\nu(x, 0) \right) = 0 \quad (12)$$

Particle set after symmetry breaking, unitary gauge

- ▶ A_μ — localized zero mode, photon.
- ▶ W_μ^\pm — resides on the brane.
- ▶ B_μ^m — continuous spectrum of massive modes.

Particle localization

Integration out compact extra dimension

$$\int dz \Pi_i \frac{d\theta_i}{2\pi R_i} \sqrt{|g|} \longrightarrow \int e^{-2kn|z|} dz \quad (13)$$

Effective action for B_μ

$$S_{eff} = \int d^4x dm \left(-\frac{1}{4} B_{\mu\nu} B_{\mu\nu} + \frac{m^2}{2} B_\mu B_\mu \right) + \int d^4x dm B_m(0) \bar{\psi} \gamma^\mu \psi B_\mu \quad (14)$$