

Hadronic Interaction Models & Cosmic Ray Composition

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“Avoid models as much as you can!”

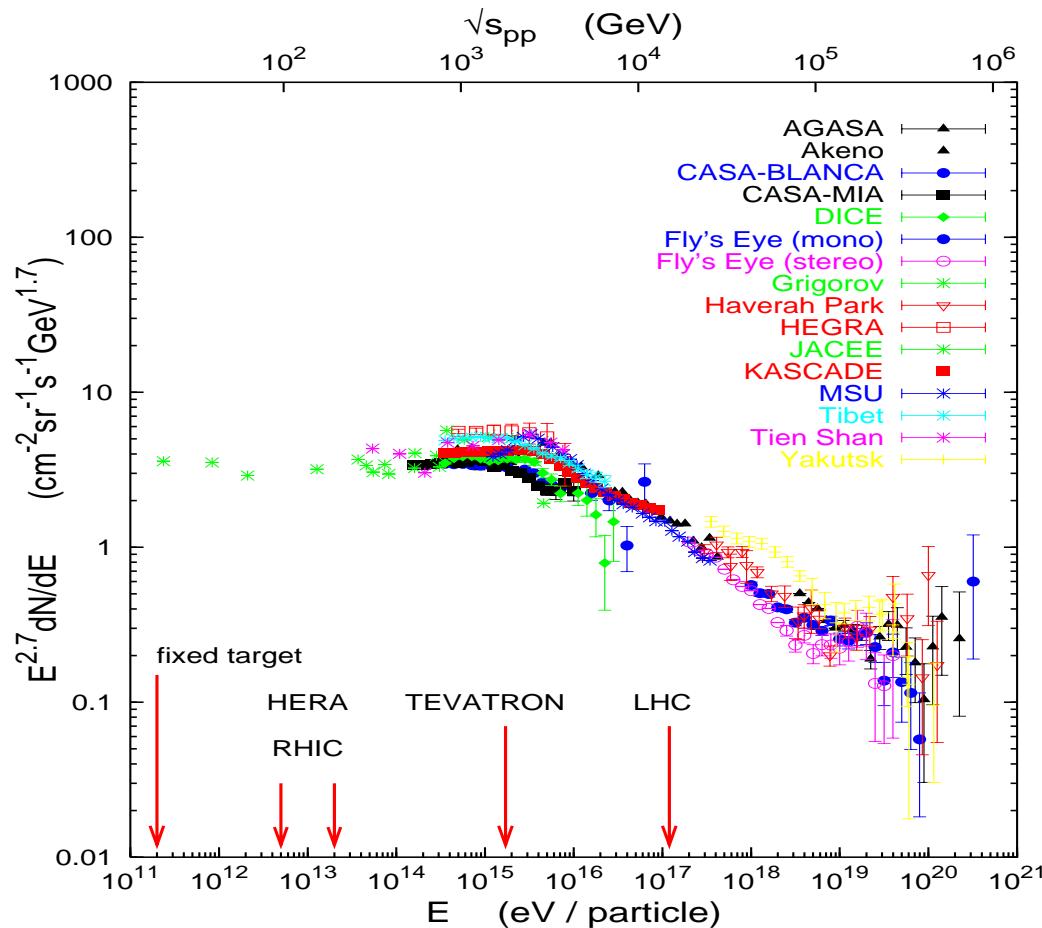
“Important issues are INPUT OF REAL DATA ...”

A. Watson

Layout

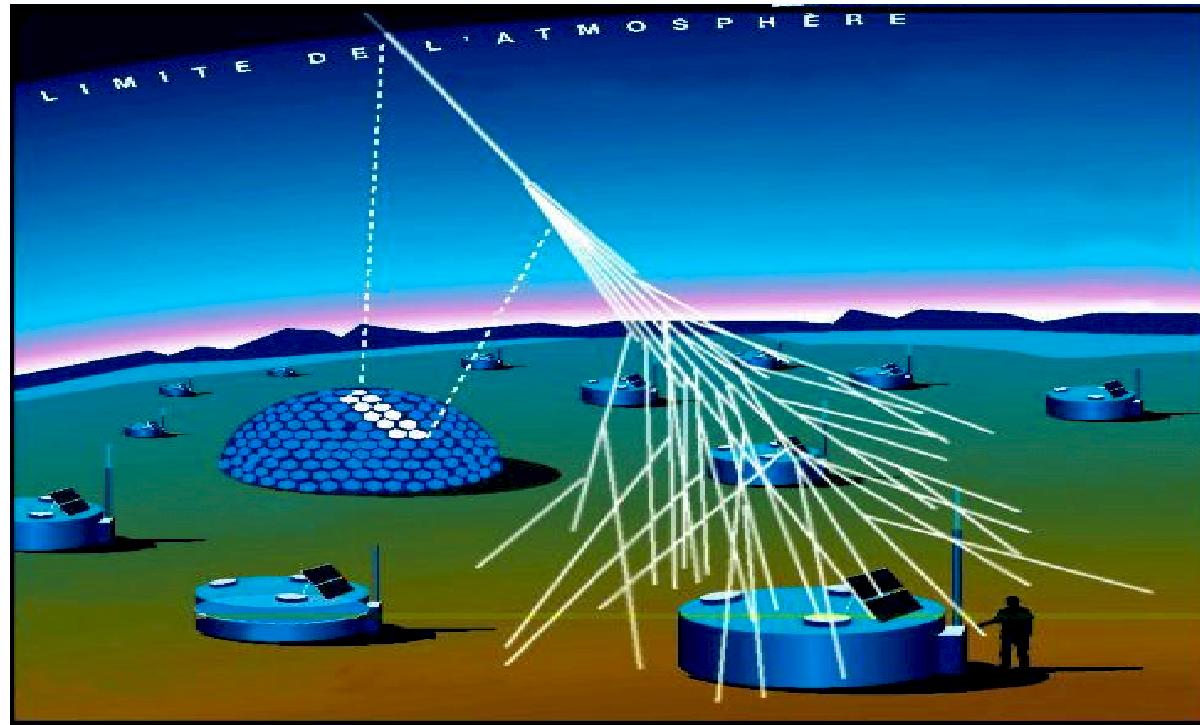
- EAS techniques & models of hadronic interactions
- High energy interactions: basic scheme
 - qualitative picture
 - 'elementary' interaction
 - multiple scattering approach
 - non-linear effects
- CR composition
 - 'knee' region & composition studies by KASCADE
 - transition to very high energies & KASCADE-Grande
 - UHECR: HiRes & Pierre Auger
- Potential changes?

In CR applications models are probed far beyond accelerator range



⇒ allows to test model consistency at higher energies & different kinematics

EAS technique of UHECR detection



- ground observations (using the atmosphere as the target)
 - primary energy \iff charged particle density at ground
 - CR composition \iff muon density at ground
- measurements of fluorescence light
 - primary energy \iff integrated light
 - CR composition \iff shower maximum position X_{\max}

Data analysis requires good understanding of air shower development

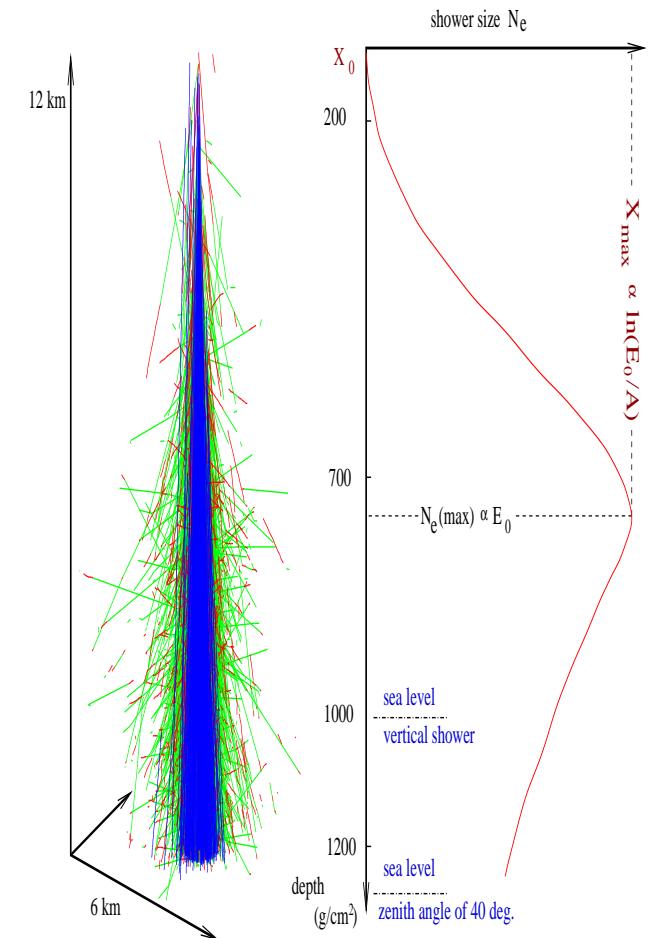
- backbone - hadron cascade
- \Rightarrow interaction models important

Requirements to CR interaction models:

- cross section predictions
- description of general hA (AA) collisions
- importance of “forward” region
- predictive power (no re-tuning possibilities)

But:

- low sensitivity to “fine” details
(smoothed by EAS development)
- charm, ..., new rare processes - irrelevant:
 - small inclusive cross sections
 - produced mainly at central rapidities



\Rightarrow CR interaction models \equiv models of “typical” (mb level) interactions

Example models:

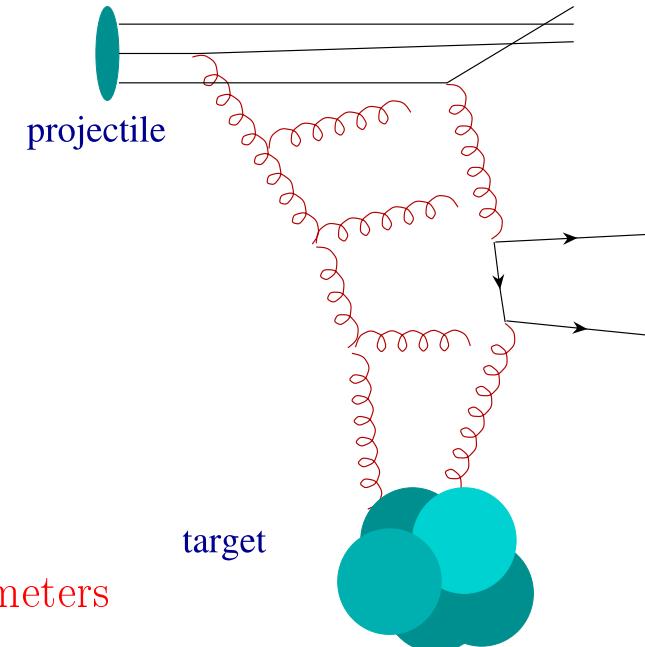
- SIBYLL 2.1 (Engel, Gaisser, Lipary & Stanev): CR analog of PITHYA
- QGSJET / QGSJET-II (Kalmykov & SO / SO): based on QGS model (Kaidalov & Ter-Martyrosian, 1982) approach; Pomeron-Pomeron interactions
- EPOS (Pierog & Werner): separate treatment of 'dense' (central 'core') and 'dilute' (peripheral 'corona') interaction regions; parameterized treatment of saturation

High energy interactions: qualitative picture

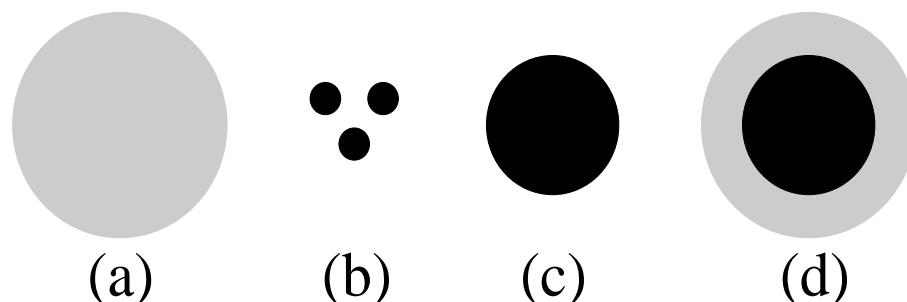
Hadronic interactions - multiple scattering processes (parton cascades):

Single scattering:

- (a) 'soft' (all $|q^2| \sim p_t^2 < Q_0^2$, $Q_0 \sim 1 \text{ GeV}^2$) cascade
 - large effective area ($\Delta b^2 \sim 1/|q^2|$)
 - slow energy rise \Rightarrow dominant at relatively low energies
- (b) cascade of 'hard' partons (all $|q^2| \gg Q_0^2$)
 - small effective area
 - rapid energy rise \Rightarrow important at very high energies and small impact parameters
- (c) 'semi-hard' scattering (some $|q^2| > Q_0^2$)
 - large effective area
 - rapid energy rise \Rightarrow dominates at high energies and over a wide b -range



(picture from R. Engel)



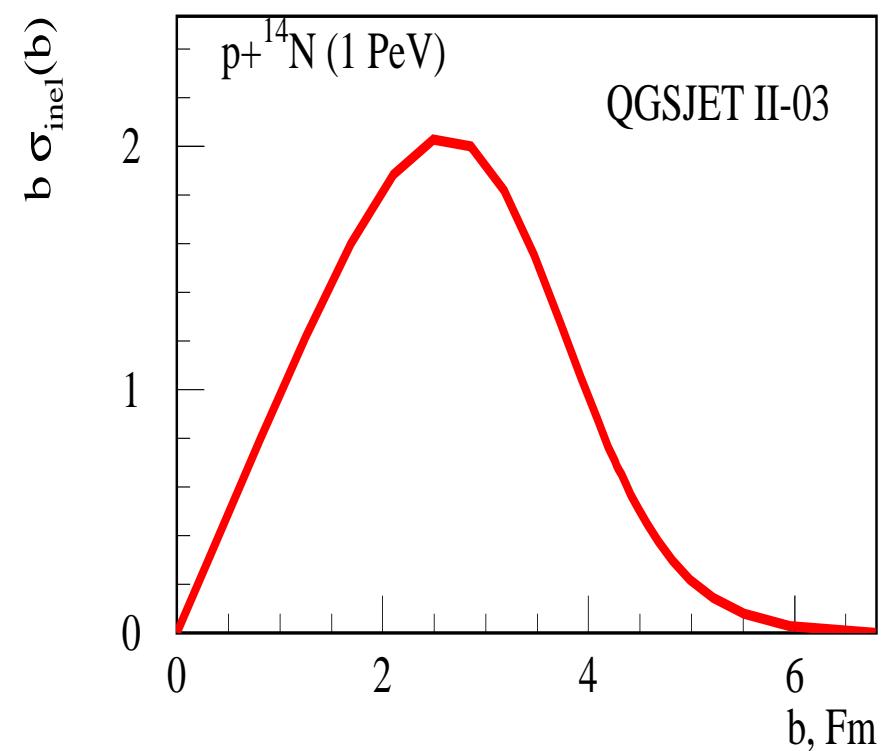
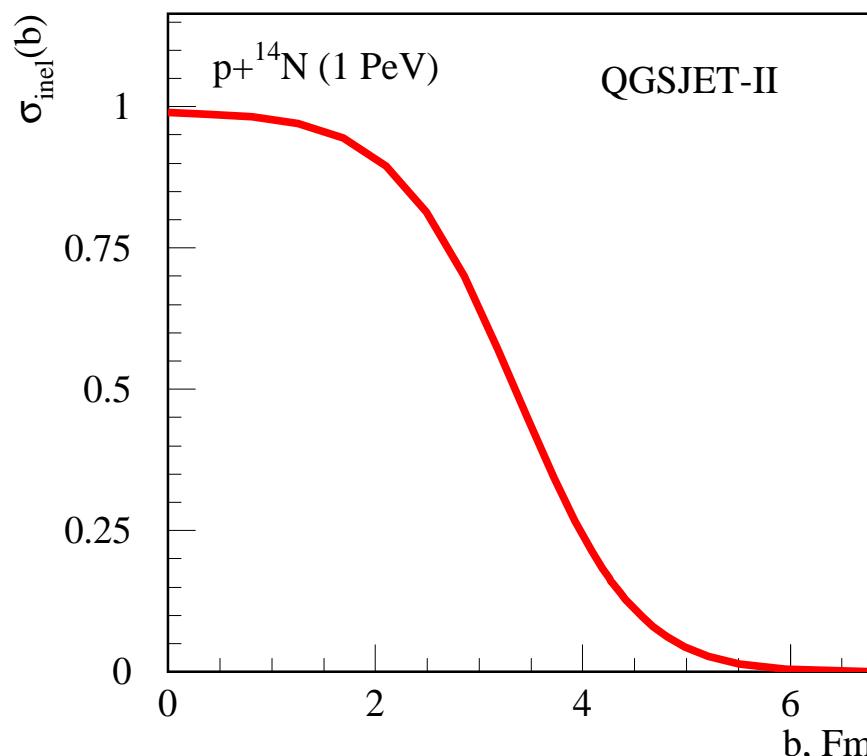
'Central' & peripheral collisions:

relative importance for $N_{h\text{-air}}^{\text{ch}}$?

What is 'central'?

- 'black disc' limit: $\sigma^{\text{inel}}(b) \sim 1 \Rightarrow \sigma^{\text{el}}/\sigma^{\text{tot}} \simeq 1/2$
- experiment: $\sigma_{pp}^{\text{el}}/\sigma_{pp}^{\text{tot}} \simeq 1/4$ @ $\sqrt{s} = 1.8$ TeV

Interaction profile & b -contributions to $\sigma_{p\text{-air}}^{\text{inel}}$ @ $E_0 = 10^6$ GeV:



'Elementary' interaction

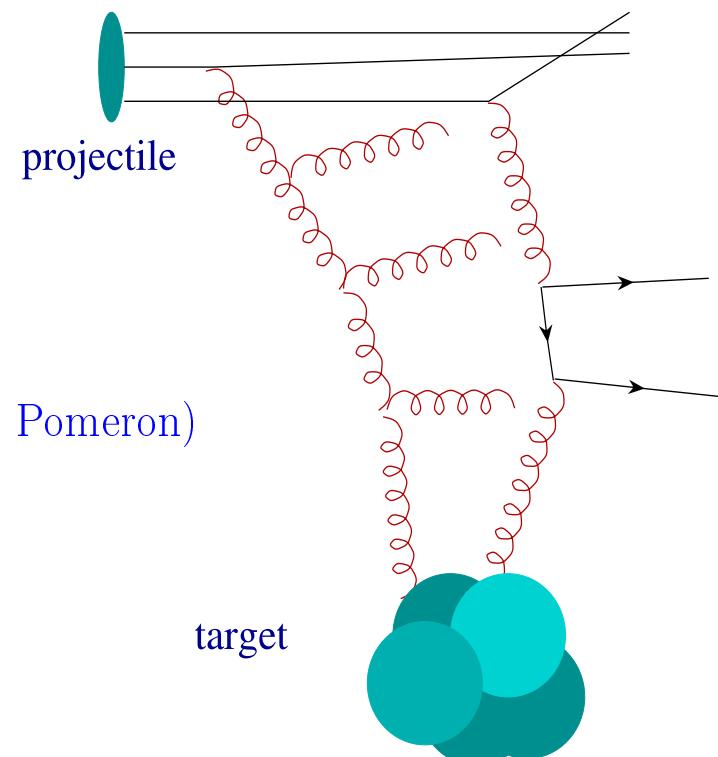
General model strategy:

- describe 'elementary' interactions (parton cascades)
 - scattering amplitude
 - hadronization procedure (conversion of partons into hadrons)
- apply Reggeon approach to treat multiple scattering processes
- describe particle production as a superposition of a number of 'elementary' processes

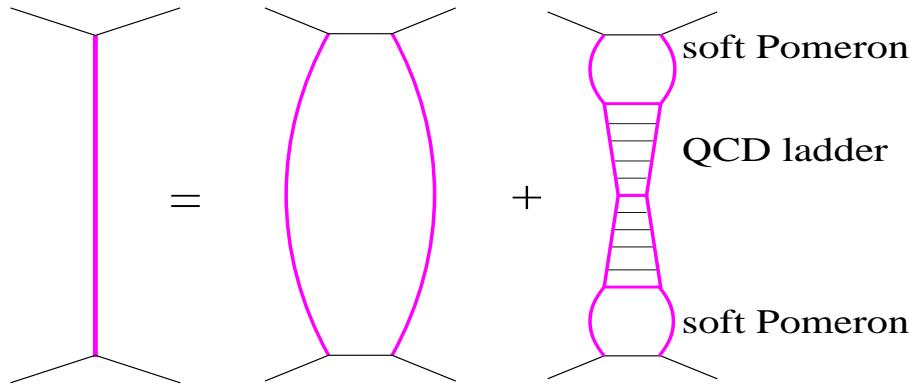
Parton cascades **start at low virtualities**

For example (QGSJET, neXus, EPOS):

- Q_0^2 - cutoff between 'soft' and perturbative physics
- 'Soft' interactions (all $|q^2|$ small $\Rightarrow \alpha_s(q^2) > 1$):
 - pQCD is inapplicable \Rightarrow Regge pole amplitude ('soft' Pomeron)
- 'Semi-hard' processes ($|q^2| > Q_0^2 \Rightarrow \alpha_s(q^2) \ll 1$)
 - 'soft' Pomeron for $|p_t^2| < Q_0^2$
 - QCD parton ladder for $|p_t^2| > Q_0^2$



General interaction \Rightarrow 'general Pomeron':



- particle production: perturbative (=calculable) parton cascade + string hadronization

'Minijet' scheme (e.g., SIBYLL) - similar but without 'soft pre-evolution' \Rightarrow smaller K_{inel} , $\langle N_{\text{ch}} \rangle$

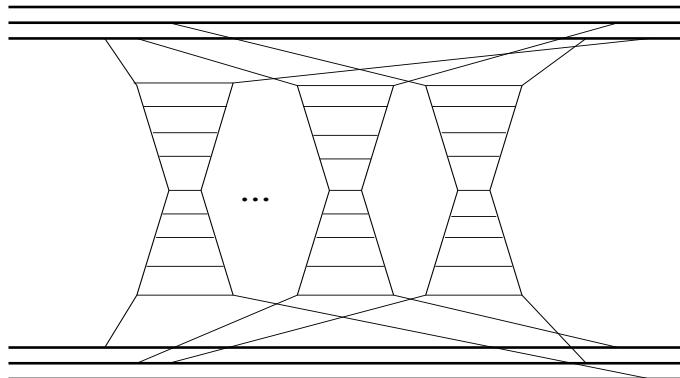
Important: direct relation between the eikonal and inclusive jet cross section

$$\begin{aligned}
 2 \int d^2 b \chi_{ad}^{\mathbb{P}_{\text{sh}}}(s, b) &\equiv \sigma_{ad}^{\text{jet}}(s, Q_0^2) \\
 \sigma_{ad}^{\text{jet}}(s, Q_0^2) &= \sum_{I,J=q,\bar{q},g} \int_{p_t^2 > Q_0^2} dp_t^2 \int dx^+ dx^- \frac{d\sigma_{IJ}^{2 \rightarrow 2}(x^+ x^- s, p_t^2)}{dp_t^2} \\
 &\times f_{I/a}(x^+, M_F^2) f_{J/d}(x^-, M_F^2)
 \end{aligned}$$

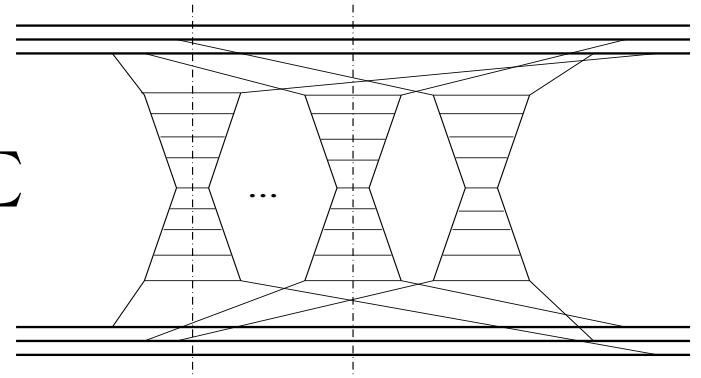
Reggeon calculus (Gribov, 1968)

General interaction - superposition of many 'elementary' processes:

$$\sigma_{ad}^{\text{tot}} = \text{Im}$$

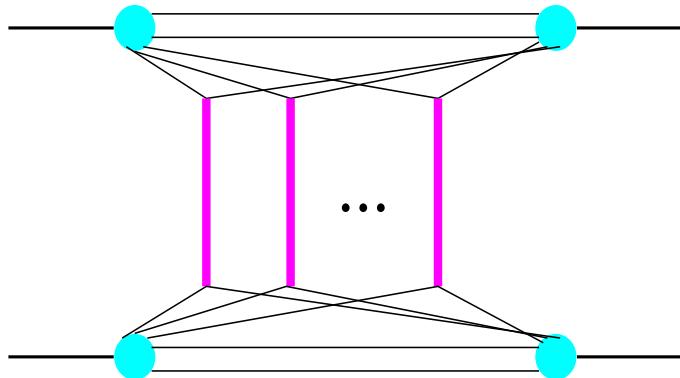


$$= \sum$$

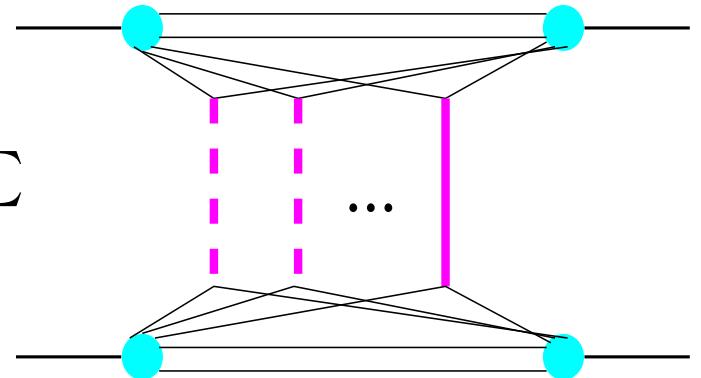


Or (Abramovskii-Gribov-Kancheli cutting rules)

$$\sigma_{ad}^{\text{tot}} = \text{Im}$$



$$= \sum$$



Contribution with m inelastic and n elastic subprocesses (diffraction neglected):

$$\sigma_{ad}^{(m,n)}(s) = \int d^2 b \frac{[2\chi_{ad}^{\mathbb{P}}(s, b)]^n}{n!} \frac{[-2\chi_{ad}^{\mathbb{P}}(s, b)]^m}{m!}$$

Physical quantity - 'topological' cross sections (m inelastic processes):

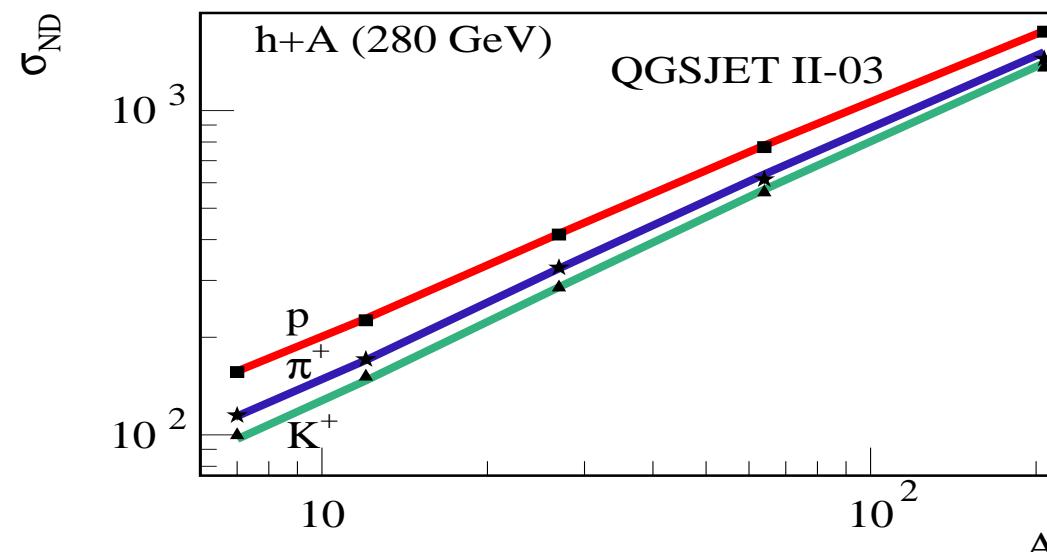
$$\sigma_{ad}^{(m)}(s) = \sum_{n=0}^{\infty} \sigma_{ad}^{(m,n)}(s) = \int d^2b \frac{[2\chi_{ad}^{\mathbb{P}}(s, b)]^n}{n!} e^{-2\chi_{ad}^{\mathbb{P}}(s, b)}$$

\Rightarrow inelastic cross section:

$$\sigma_{ad}^{\text{inel}}(s) = \sum_{m=1}^{\infty} \sigma_{ad}^{(m)}(s) = \int d^2b \left[1 - e^{-2\chi_{ad}^{\mathbb{P}}(s, b)} \right]$$

Parameter-free generalization to hadron-nucleus (nucleus-nucleus) scattering

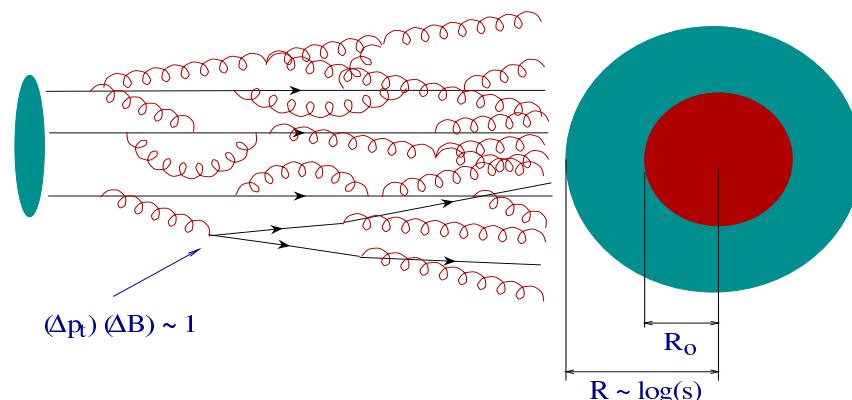
E.g., predicted non-diffractive hadron-nucleus cross section vrs. experimental data



Non-linear effects

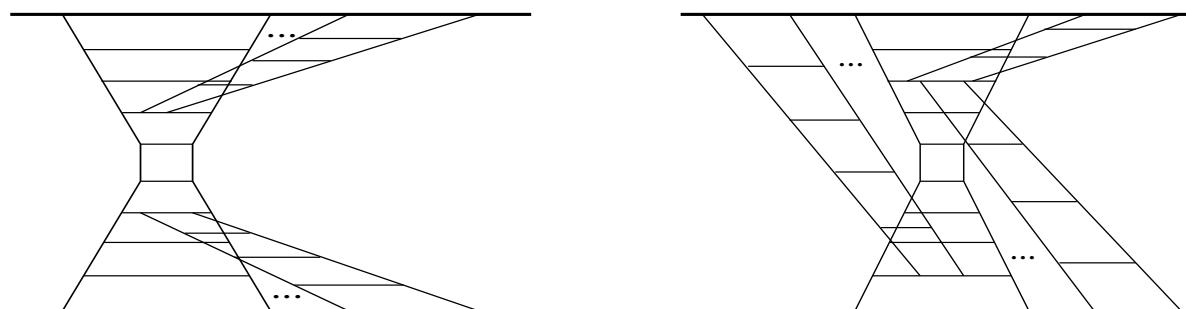
Large s , small b , large A :

- many partons closely packed
- \Rightarrow parton cascades overlap and interact with each other
- \Rightarrow parton shadowing (slower rise of parton density)
- saturation (maximal possible density reached)

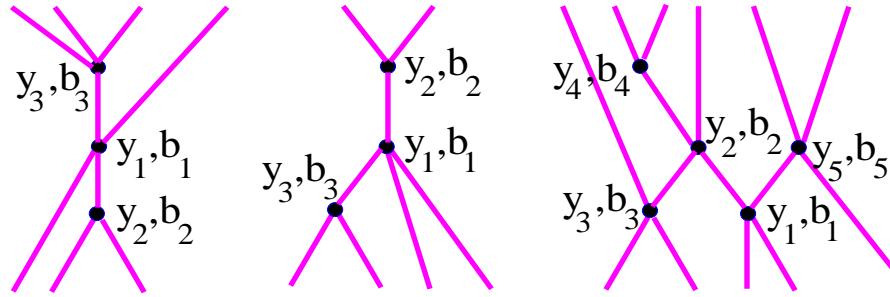


(picture from R. Engel)

Non-linear effects in QCD: interaction between parton ladders



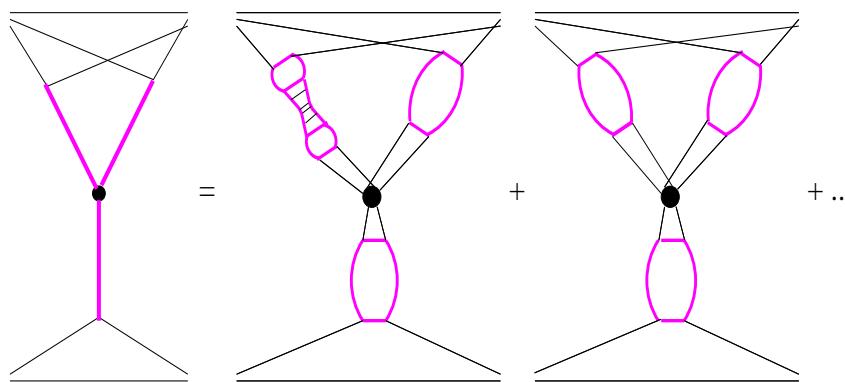
Pomeron approach: non-linear effects \equiv Pomeron-Pomeron interactions



QGSJET II: all order re-summation of Pomeron 'nets' (SO, 2006, 2008)

Basic assumptions:

- neglect saturation effects above a **fixed** Q_0^2 scale
- Pomeron-Pomeron coupling - only at $|q|^2 < Q_0^2$
- \Rightarrow only "soft" Pomeron coupling



Summary contribution of all 'net'-like enhanced graphs:

$$\sum_{\substack{m_1, n_1 \geq 1 \\ m_1 + n_1 \geq 3}} - \sum_{\substack{m_1, n_2 \geq 1 \\ m_2, n_1 \geq 0 \\ m_i + n_i \geq 2}} \text{---}$$

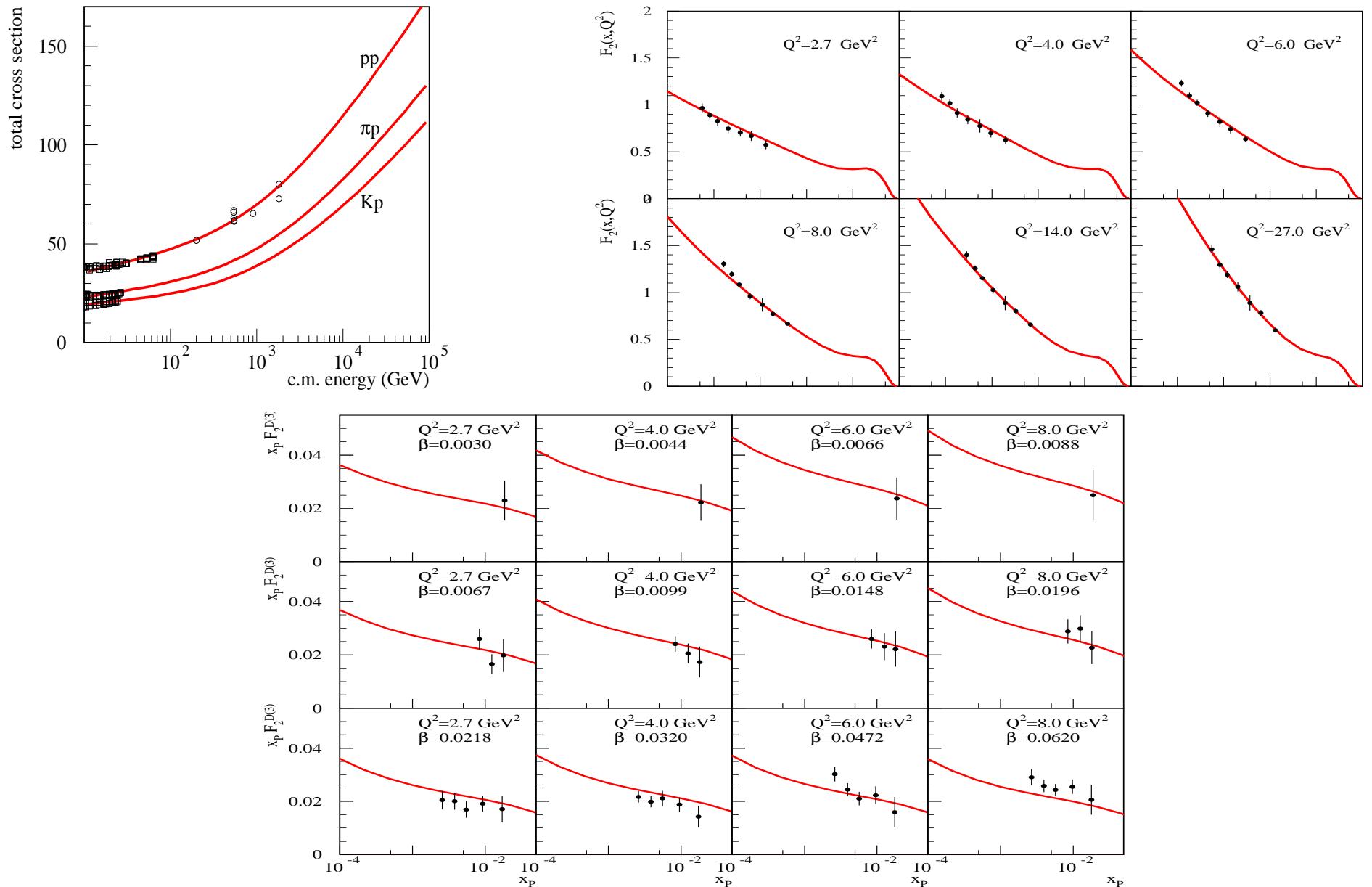
Expressed via 'net-fan' contributions (parton (y, b) -distributions 'probed' during interaction):

$$\text{net} = \text{---} + \sum_{\substack{m_2 \geq 1, n_2 \geq 0 \\ m_2 + n_2 \geq 2}} \text{---}$$

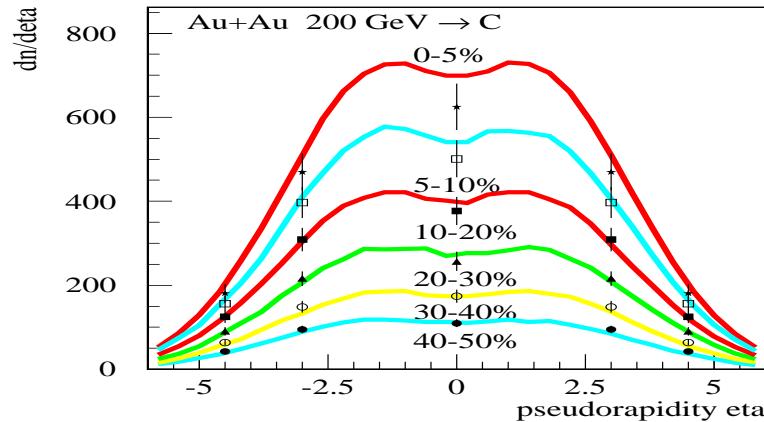
Similar (recursive) representations obtained for the complete set of AGK cuts

\Rightarrow MC implementation

- allows to get consistency between cross sections & structure functions (total and diffractive):



- (parameter-free) A -enhancement of screening effects in hA (AA) collisions (e.g., at RHIC)



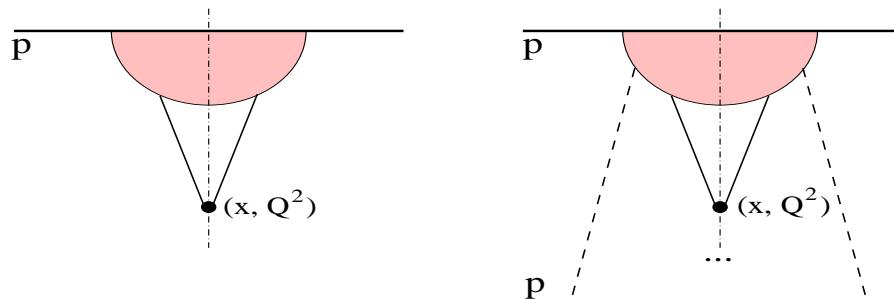
- preserves QCD factorization for inclusive jet cross section

$$\sigma_{ad}^{\text{jet}}(s, Q_0^2) = \sum_{I,J=q,\bar{q},g} \int_{p_t^2 > Q_0^2} dp_t^2 \int dx^+ dx^- f_{I/a}^{\text{scr}}(x^+, M_F^2) f_{J/d}^{\text{scr}}(x^-, M_F^2) \frac{d\sigma_{IJ}^{2 \rightarrow 2}(x^+ x^- s, p_t^2)}{dp_t^2}$$

- but: $\int d^2b \chi_{ad}^{\mathbb{P}_{\text{sh}}}(s, b) \neq \sigma_{ad}^{\text{jet}}(s, Q_0^2)$ - due to non-factorizable contributions

Why?

- inclusive spectra - expressed via PDFs of free hadrons
- eikonal - expressed via 'parton distributions' probed during interaction:

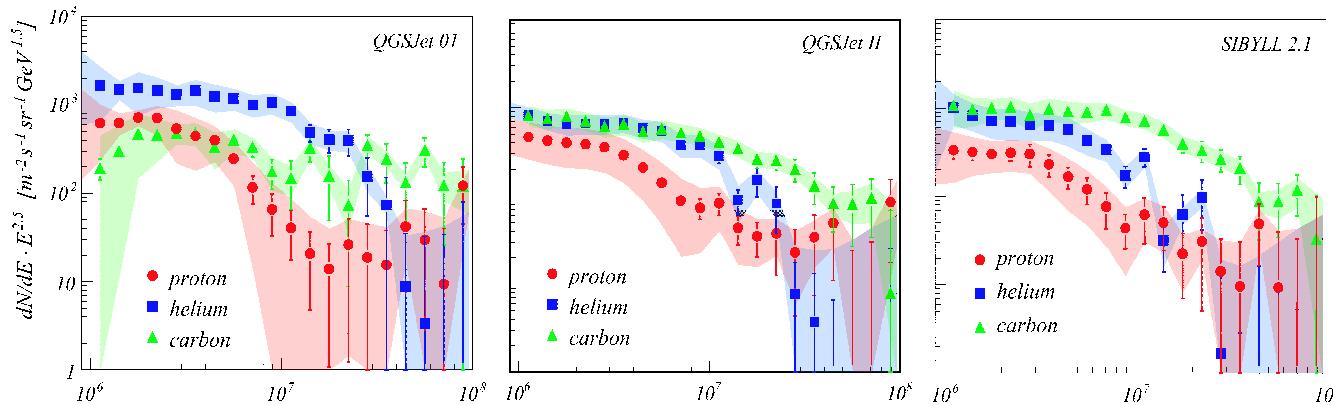


CR composition

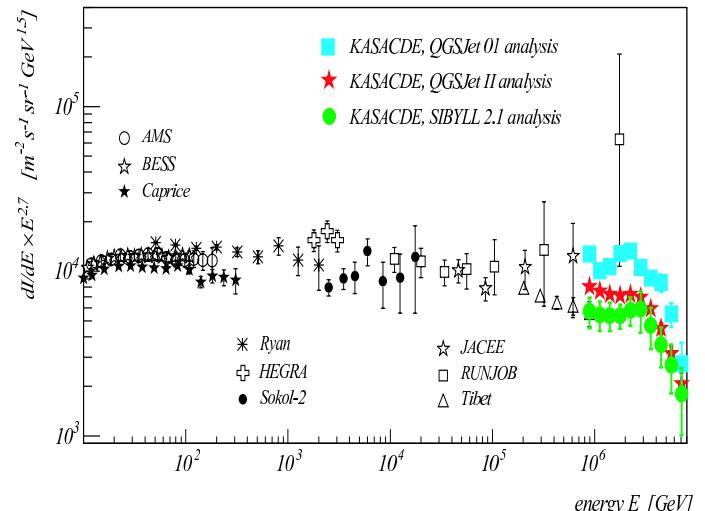
CR studies in the 'knee' region ($E_0 \sim 10^{15} \div 10^{16}$ eV) - generally consistent with each other

KASCADE N_e/N_μ -deconvolution:

- a moderate model dependence
- consistent with rigidity-dependent 'knee' position (at least, for p & He)

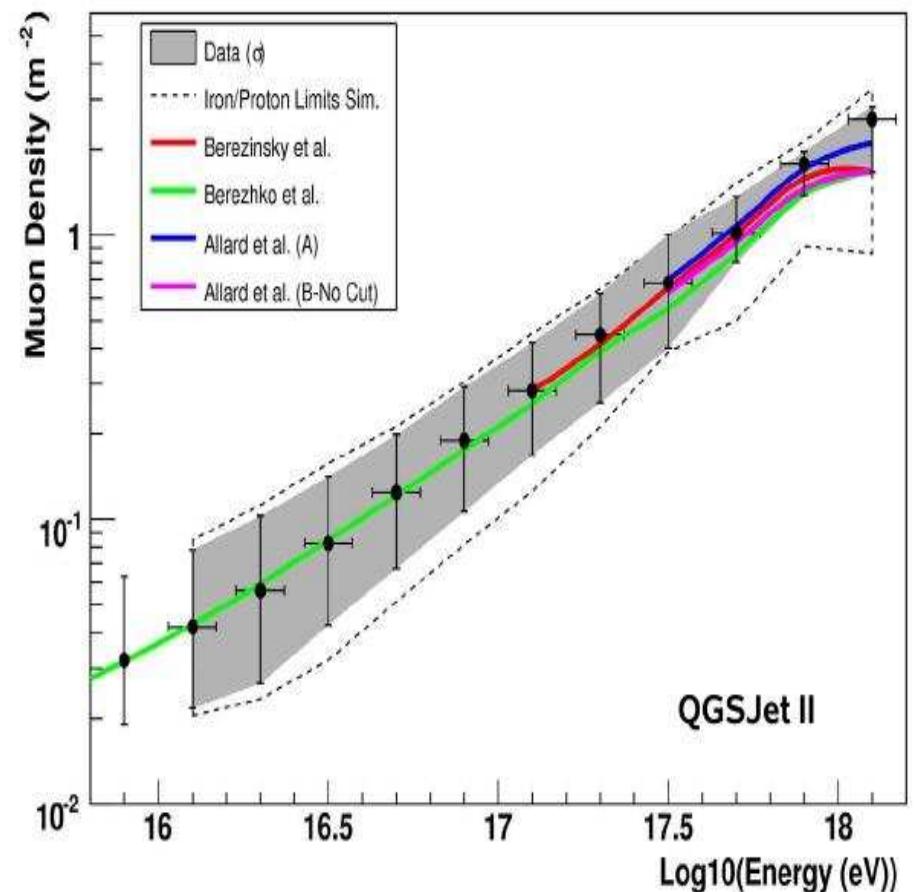
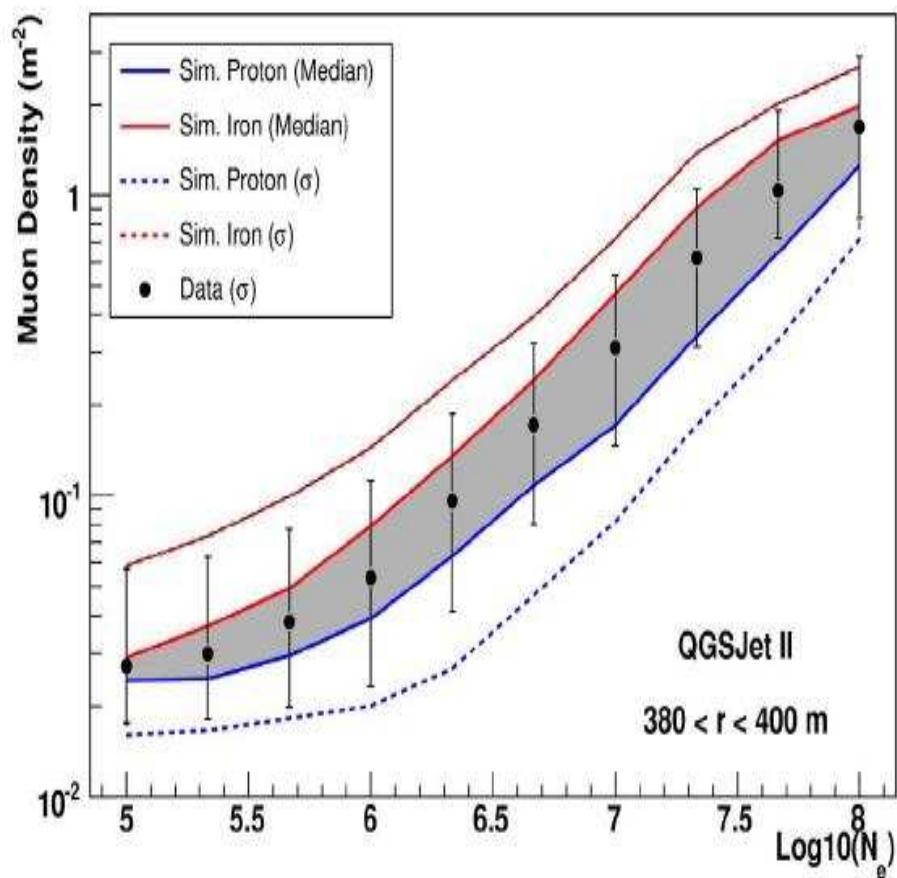


Proton spectrum - consistent with direct measurements:

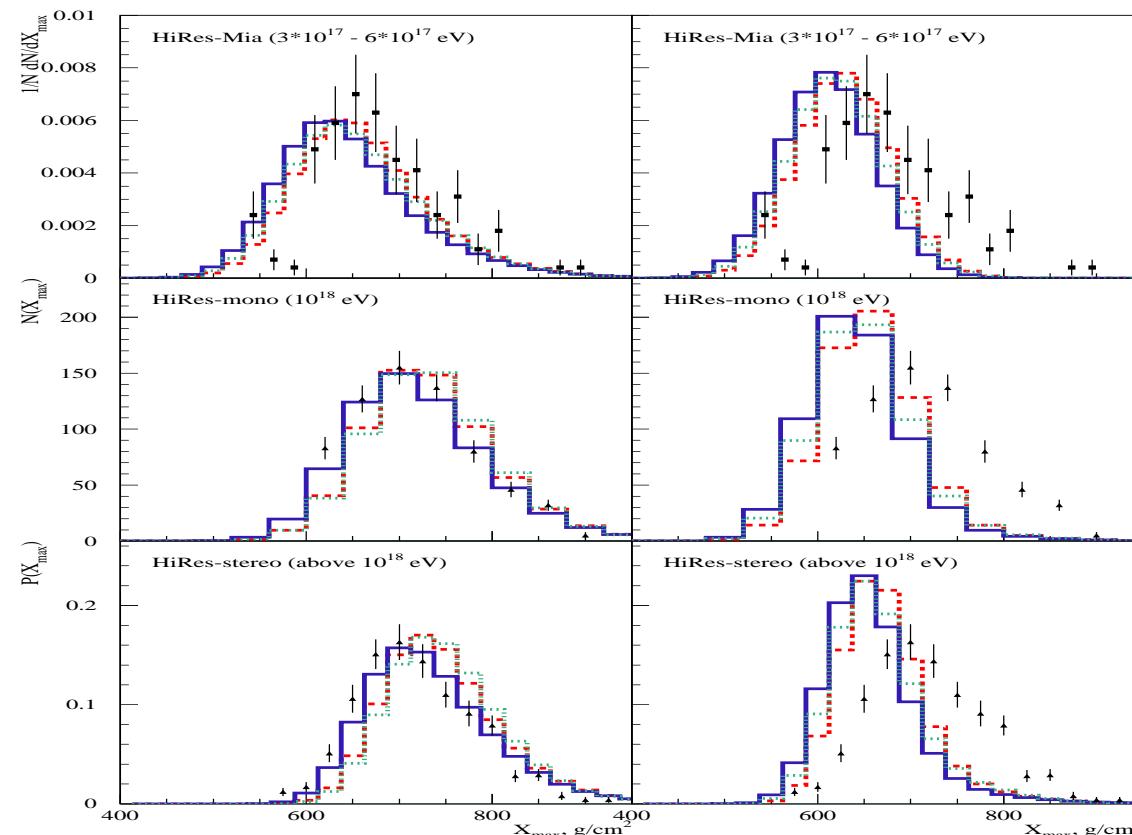
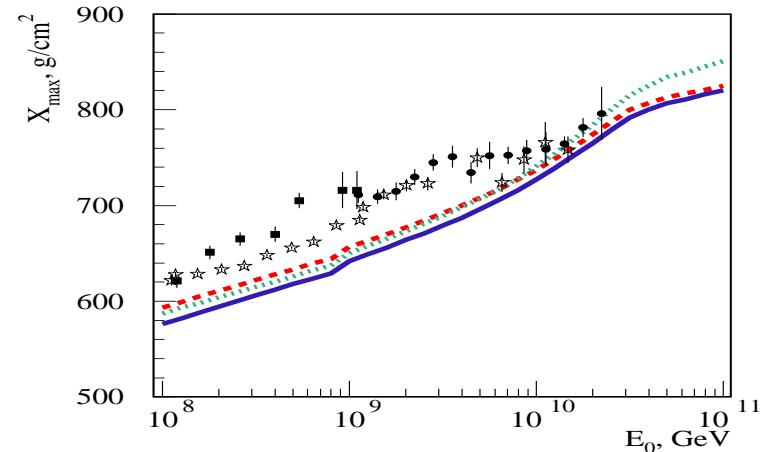
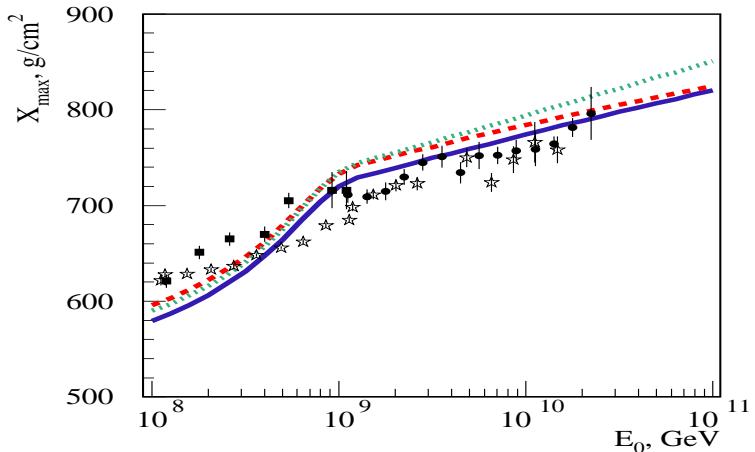


KASCADE-Grande: $\rho_\mu(400 \text{ m})$ up to 10^{18} eV

- consistent with model predictions (bounded by p / Fe)
- consistent with astrophysical models (but can not discriminate them)

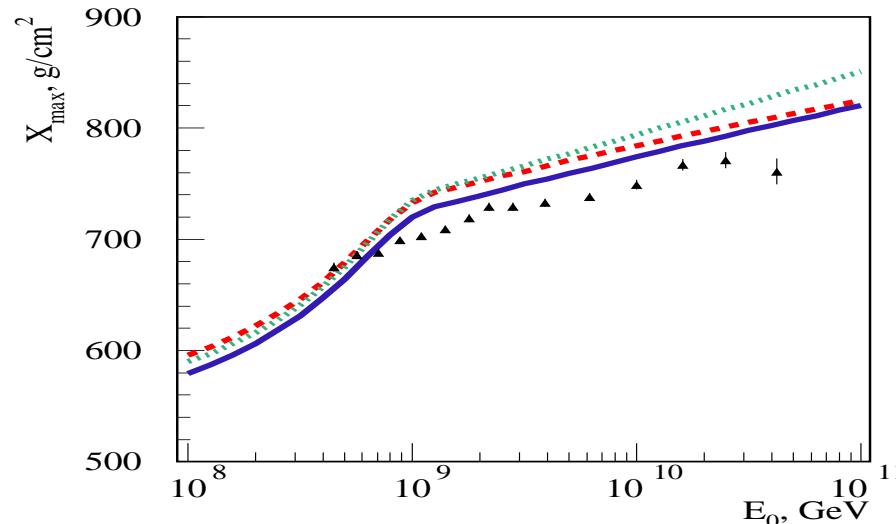


Galactic-extragalactic transition: HiRes data rather favor the 'dip' scenario, reject 'ankle' models

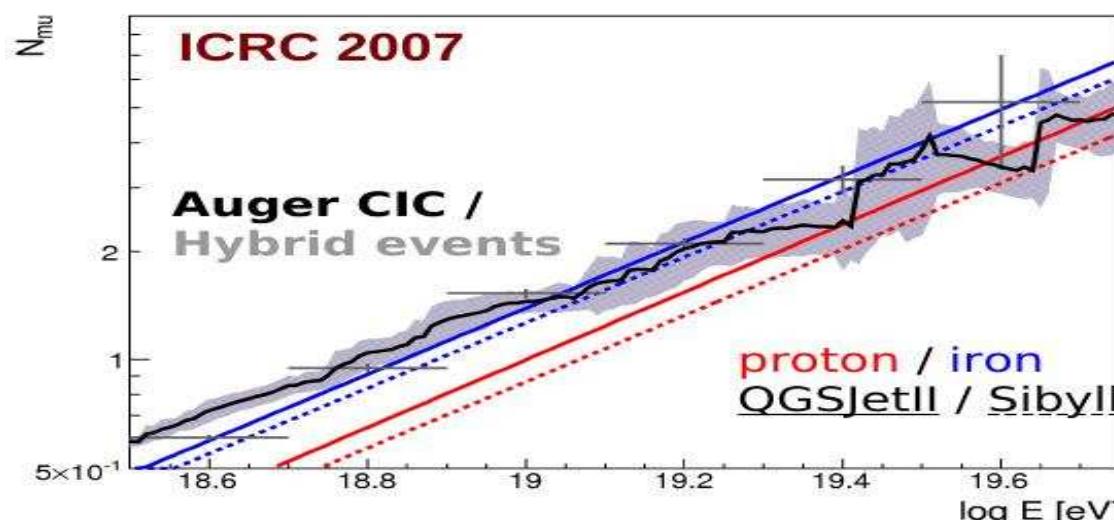


Pierre Auger AGN correlations - also support the 'dip' picture

Contradiction with Pierre Auger composition studies:



Strong contradiction with Pierre Auger results on N_μ (based on EAS universality, CIC):



Potential changes?

QGSJET-II update: 'Pomeron loops' (in progress):

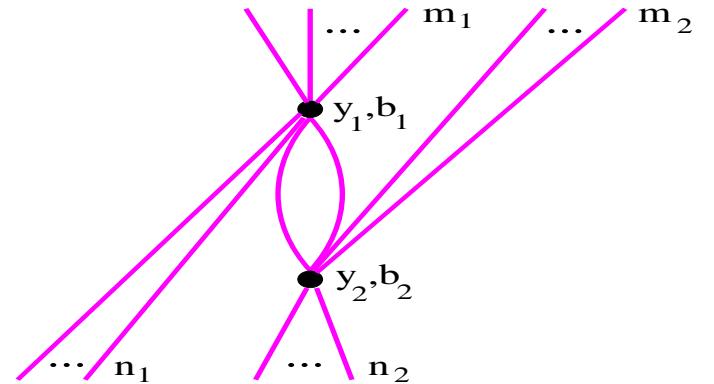
- small at low parton density ($\sim G^2$)
- suppressed at high density:

$$\sim \sum_{n_1=0}^{\infty} \frac{(-\chi_{d\P}^{\mathbb{P}}(s_0 e^{y_1}, b_1))^{n_1}}{n_1!} = e^{-\chi_{d\P}^{\mathbb{P}}(s_0 e^{y_1}, b_1)}$$

Still a finite correction at large b :

- of importance for σ^{tot} , σ^{diffr}
- will lead to smaller 3P-coupling
- \Rightarrow smaller nuclear screening
- \Rightarrow higher $N_{\text{h-air}}^{\text{ch}}$, $K_{\text{h-air}}^{\text{inel}}$

But only a **moderate effect** expected



What else:

- BFKL evolution: at NLO close DGLAP
- saturation at $Q^2 > Q_0^2$: irrelevant for σ_{inel} \Rightarrow additional suppression of particle production
- independent hadronization of valence quarks: for central collisions \Rightarrow hard to expect a big effect

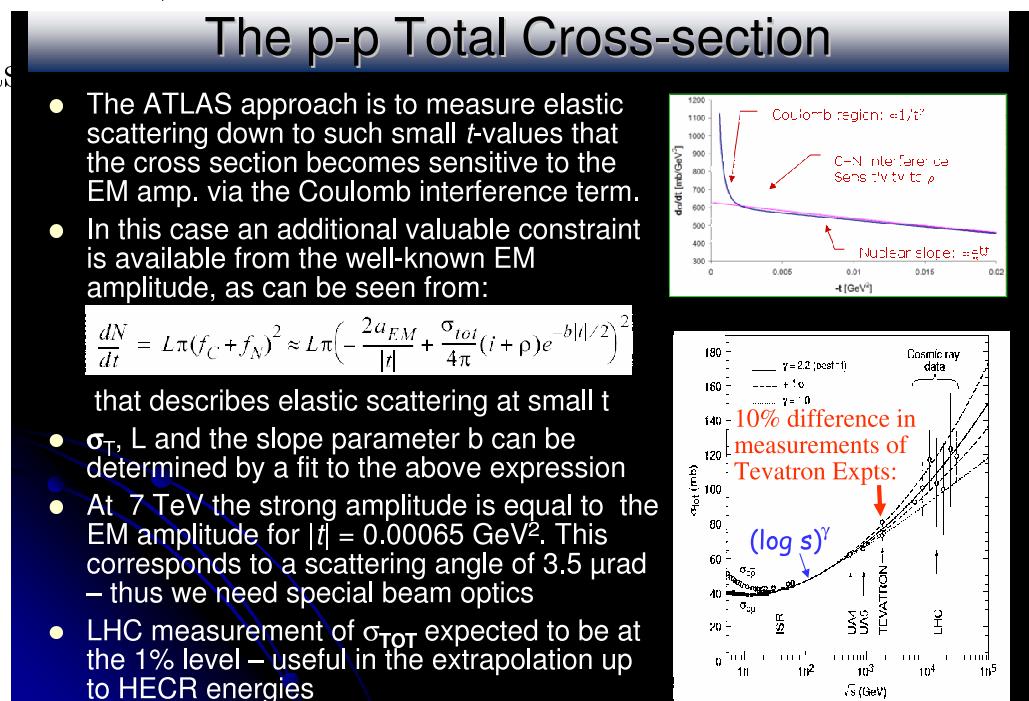
Extra slides

Hadronic cross sections - of crucial importance for EAS applications

- $\sigma_{h\text{-air}}^{\text{inel}}$ - direct impact on X_{\max}
- model calibration:
 - particle production: **mainly** with fixed target hp data
 - energy extrapolation: **mainly** inferred from $\sigma_{pp}^{\text{tot}}(s)$ behavior

⇒ measurement of σ_{pp}^{tot} with 1% accuracy (~ 10 mb) - most important LHC contribution:

- allows to obtain $\sigma_{p\text{-air}}^{\text{inel}}$ (Glauber + inelastic)
- significantly improves model calibration



FP420 experiment:

- designed to study diffractive Higgs production
- can measure 'soft' diffraction

The FP420 Project

$M^2 = z_1 z_2 S$
Where $z_{1,2}$ are the fractional momentum losses of the outgoing protons

FP420 TOTEM / ATLAS RPs Central Detector System
Leading proton CMS/ATLAS Leading proton
420 m 308 m 215 m

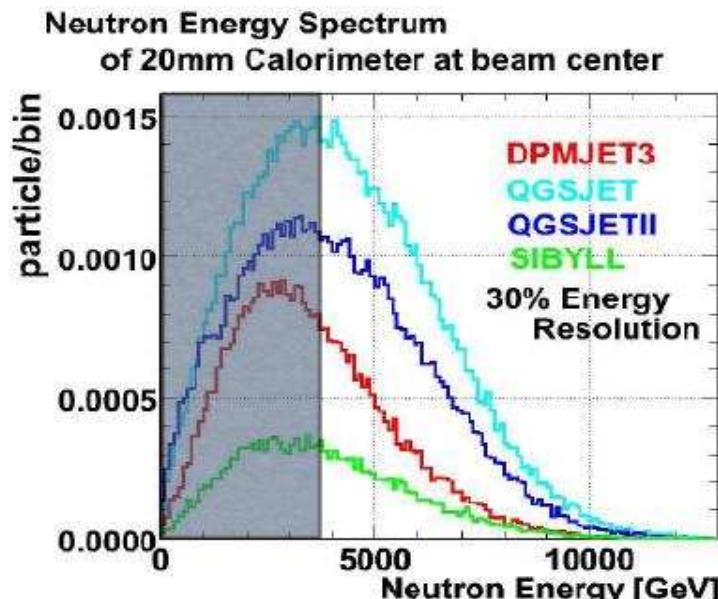
- Extend the acceptance for leading proton tagging
- Combine information from central detector and RP @220m

- Exclusive central Higgs prod. $pp \rightarrow p H p$: 3-10fb
- Inclusive central Higgs prod. $pp \rightarrow p+X+H+Y+p$: 50-200 fb
- Reconstruction of the central mass:
$$M_H^2 = (p + \bar{p} - p' - \bar{p}')^2 \quad \Delta M = O(1.0 - 2.0) \text{ GeV}$$

FP420: R&D fully funded

- TDR to ATLAS/CMS by 1st half of '07 then to the LHC.

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LHCf Experiment

Low energy beam profile at the top of the detector
beam pipe detector box TAN area interaction point
96mm 140m away 140m

1 TeV neutron 400 GeV photon

- LHCf: measurement of photons and neutral pions in the very forward region of LHC
- Add an EM calorimeter at 140 m from the Interaction Point (of ATLAS) (Scintillating fiber /Tungsten calorimeter + Silicon strip det. Calorimeter)
- At the LHC the 14 TeV Ecm translates to a 10^{17} eV Lab. Energy - by comparing experimental results with MC predictions one can tune MC used in cosmic ray EAS simulation.

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Good-Walker scheme for diffraction

Low mass diffraction - 2-component eikonal scheme: $|a\rangle = \frac{1}{\sqrt{2}}|1_a\rangle + \frac{1}{\sqrt{2}}|2_a\rangle$

$|k_a\rangle$ - diffractive eigenstates for hadron a ;

couple to a Pomeron with different strength $\lambda_{k/a}$ ($\lambda_{k/a} + \lambda_{k/a} = 2$)

$$\sigma_{pp}^{\text{inel}}(s) = \int d^2b \frac{1}{4} \sum_{i,j=1}^2 \left[1 - e^{-2\lambda_{i/p}\lambda_{j/p}} \chi_{pp}^{\mathbb{P}}(s,b) \right]$$

$$\sigma_{pp}^{\text{s-diffr}}(s) = 2 \int d^2b \frac{1}{4} \sum_{j=1}^2 \left[e^{-\lambda_{1/p}\lambda_{j/p}} \chi_{pp}^{\mathbb{P}}(s,b) - e^{-\lambda_{2/p}\lambda_{j/p}} \chi_{pp}^{\mathbb{P}}(s,b) \right]^2$$

$$\frac{1}{4} \sum_{i,j=1}^2 \left[1 - e^{-2\lambda_{i/p}\lambda_{j/p}} \chi_{pp}^{\mathbb{P}}(s,b) \right] = \begin{cases} 1, & b \rightarrow 0 \\ 2\chi_{pp}^{\mathbb{P}}(s,b) - 2(2 - \lambda_{1/p}\lambda_{2/p})^2 (\chi_{pp}^{\mathbb{P}}(s,b))^2 + \dots, & b \rightarrow \infty \end{cases}$$

Higher diffraction \Rightarrow cross sections reduced at high energies (inelastic screening)

Topological cross section (n inelastic sub-processes):

$$\sigma_{ad}^{(n)}(s) = \int d^2b \frac{1}{4} \sum_{i,j=1}^2 \frac{[2\lambda_{i/a}\lambda_{j/d}\chi_{ad}(s,b)]^n}{n!} e^{-2\lambda_{i/a}\lambda_{j/d}\chi_{ad}(s,b)}$$

\Rightarrow larger multiplicity fluctuations in non-diffractive collisions

Gluon density saturation - only at small b :

