



Entropy Growth and the Equation of State

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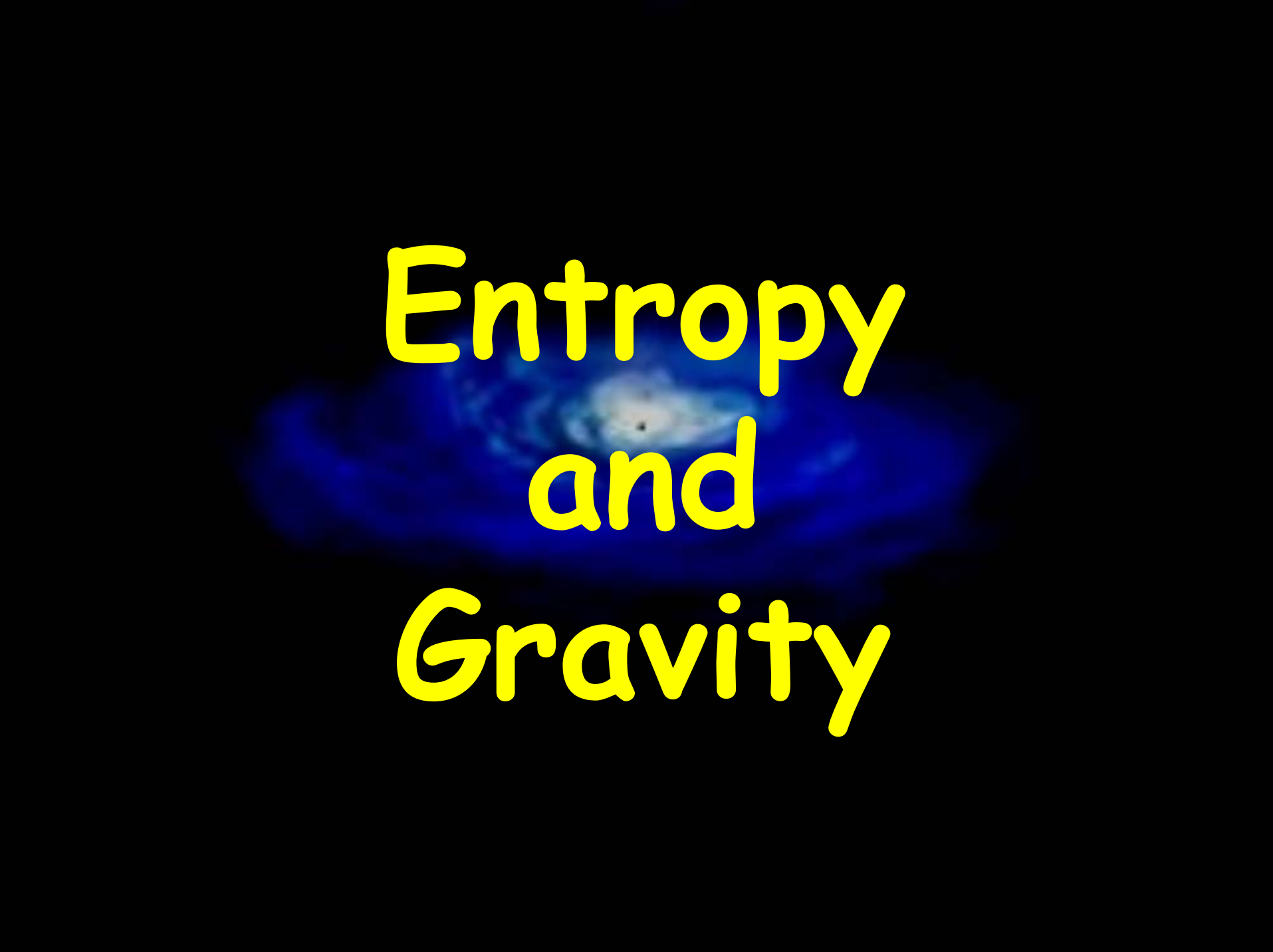
Wilfried Buchmueller

DESY

0. Overview



- 1. Entropy and Gravity
 - 2. The Second Law in Cosmological Situations
 - 3. Constraints on Dark Energy
 - 4. Summary and Conclusions
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A blue-toned image of a spiral galaxy, likely the Andromeda Galaxy, is centered in the background. The text "Entropy and Gravity" is overlaid in a bright yellow, bold, sans-serif font. The word "Entropy" is at the top, "and" is in the middle, and "Gravity" is at the bottom.

Entropy and Gravity

Black Holes and Entropy



- Black holes are rather dumb:
Their memory is limited to **mass, charge and angular momentum.**
- They do not remember the details of their origin!



- Mass is concentrated in a point.

 Seems very ordered!

 0 Entropy???



A violation of the 2. Law?



- According to the
2. Law of Thermodynamics:
Entropy always increases!
 - Throw entropic matter into BH
 - ➔ Vanishes
 - ➔ Violation of the 2. Law????
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Area Law



- The horizon area of black holes always increases!
(in classical gravitational systems)
- E.g. when matter is dropped into the black hole.

 Entropy of BH \sim Horizon Area

$$S_{\text{BH}} = 2\pi A_{\text{BH}} M_{\text{P}}^2$$

Generalized 2. Law



- In systems with ordinary matter and black holes

Generalized 2. Law (GSL)

$$dS = dS_{\text{mat}} + dS_{\text{H}} \geq 0$$

A Note on Entropy Bounds



- The GSL cannot work if an arbitrary amount of entropy can be stored in a system with fixed size and energy.

→ $S_{\text{sys}} \leq 2\pi A_{\text{sys}}$

Covariant entropy bound

$$S_{\text{sys}} \leq 2\pi E_{\text{sys}} L_{\text{sys}}$$

Bekenstein bound

Cosmological



Situations

Problems



- Space is not asymptotically flat.
 - Typically non stationary situations!
 - Global quantities like total energy and entropy are hard to define!
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Which Horizon? Which Region?



- „Horizons“:
 - Future event horizon
 - Apparent/Hubble horizon
 - Boundary of the causal region
- Regions:
 - Spacelike
 - Null „from the middle“
 - Null „from the surface“

An Example



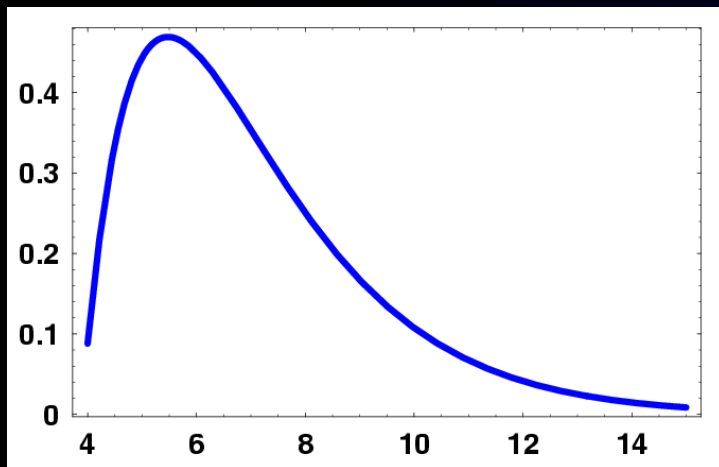
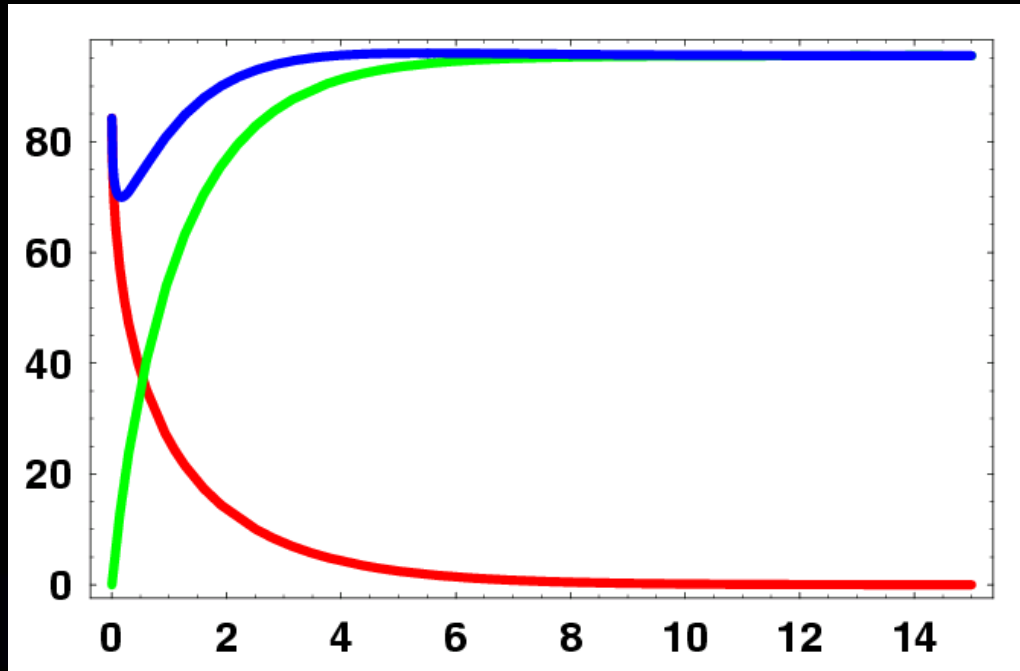
- Flat universe with Λ and radiation

$$S_{\text{H}} = 2\pi A_{\text{H}}(t) = 8\pi^2 D_{\text{H}}^2(t).$$

$$S_{\text{R}} = \frac{4}{3}\pi D_{\text{H}}^3(t) s_{\text{R}}(t) = \frac{16}{9}\pi\sigma T_0^3 \frac{D_{\text{H}}^3(t)}{a^3(t)}.$$

$$D_{\text{H}}(t) = a(t) \int_t^{\infty} \frac{dt'}{a(t')}.$$

An Example



An Example



Message:
GSL seems to
work!



**Constraints
on
Dark Energy**

Idea: (pure dark energy universe)



- Roughly speaking:

$$D_H(t) = a(t) \int_t^{\infty} \frac{dt'}{a(t')} \sim \frac{1}{H}$$

- Use Friedmann equation $3H^2 = \rho$

➡ If ρ increases

➡ D_H decreases

➡ S_H decreases ➡ Trouble!

And the Equation of State...



- Friedmann equation II:

$$\dot{\rho} = -3H\rho(1+w)$$

➔ $\dot{\rho} > 0$ for $w < -1$

➔ **Phantom Universes
may be in trouble!!**

Real Universe



- There is additional matter and radiation!

$$\rho_{\text{tot}} = \rho_{\text{DE}} + \rho_{\text{mat}} + \rho_{\text{rad}}$$

$$w_{\text{eff}} = -1 - \frac{\dot{\rho}_{\text{tot}}}{3H\rho_{\text{tot}}}$$

- Matter and radiation carry additional entropy!
- But $S_{\text{mat}} + S_{\text{rad}} \sim 10^{88+11} \ll \ll S_H \sim 10^{120}$



$S_{\text{mat}} + S_{\text{rad}}$ negligible

Assumption

- For the moment we assume that the dark energy does not carry entropy!
 - This seems reasonable:
 - A homogeneous scalar field has zero entropy
 - An adiabatically expanding perfect fluid has $\frac{S}{V} \sim \frac{1}{a^3(t)}$
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Constant equation of state



- For DE with constant EOS $w < -1/3$ we will have at late times

$$\rho_{\text{DE}} \gg \rho_{\text{mat}} + \rho_{\text{rad}}$$

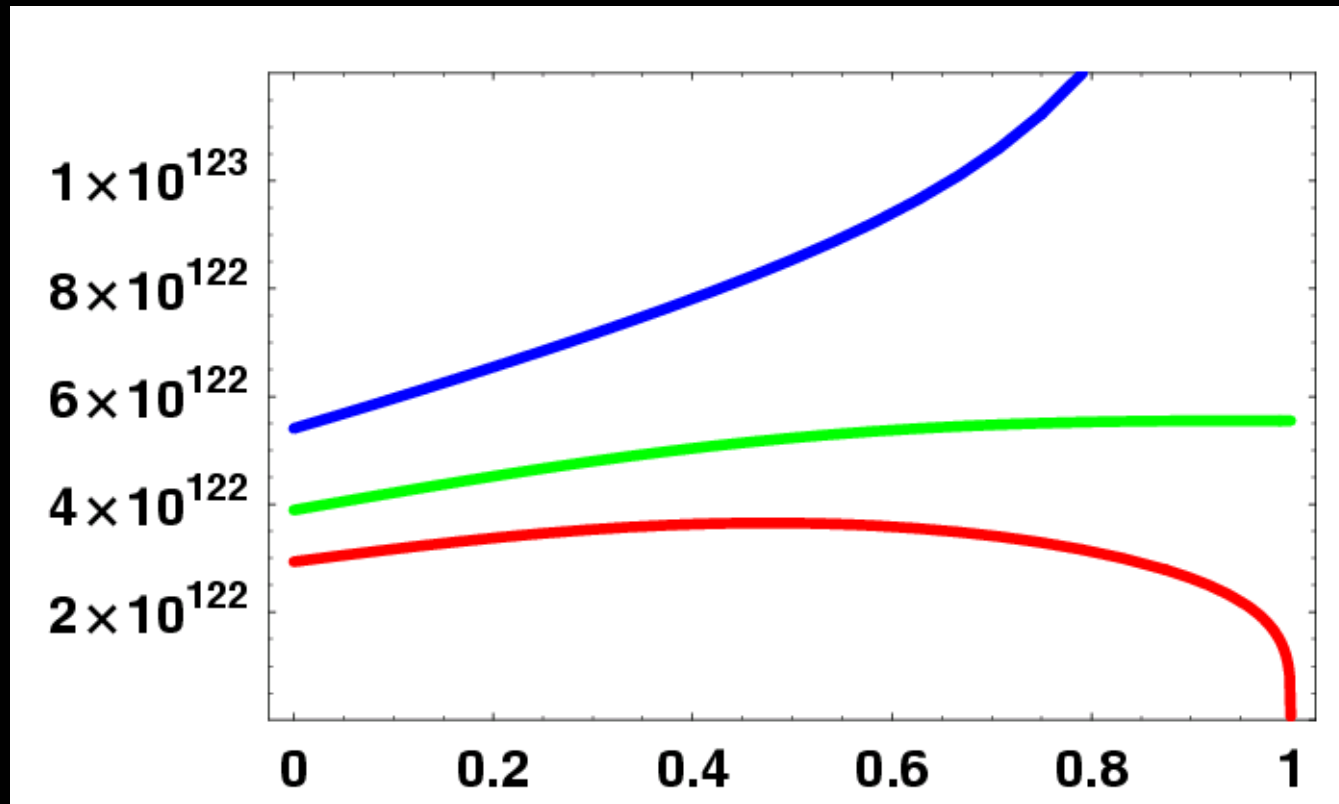
- Then $w_{\text{eff}} \approx w_{\text{DE}}$

and
$$\dot{H} = \left(\frac{\text{const}}{H} \right) \sim (1 + w_{\text{DE}})$$

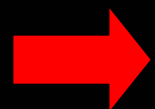
$\Rightarrow \dot{S}_{\text{H}} \geq 0 \iff w \geq -1$

\Rightarrow Phantom universes with constant EOS will violate GSL!

No Big Rip



Near Big Rip $S_{\text{tot}} \rightarrow 0$



Big Rip violates *GSL*!

Commercials

Your Universe is constantly
loosing Entropy?



Your Universe ends in a Big Rip?

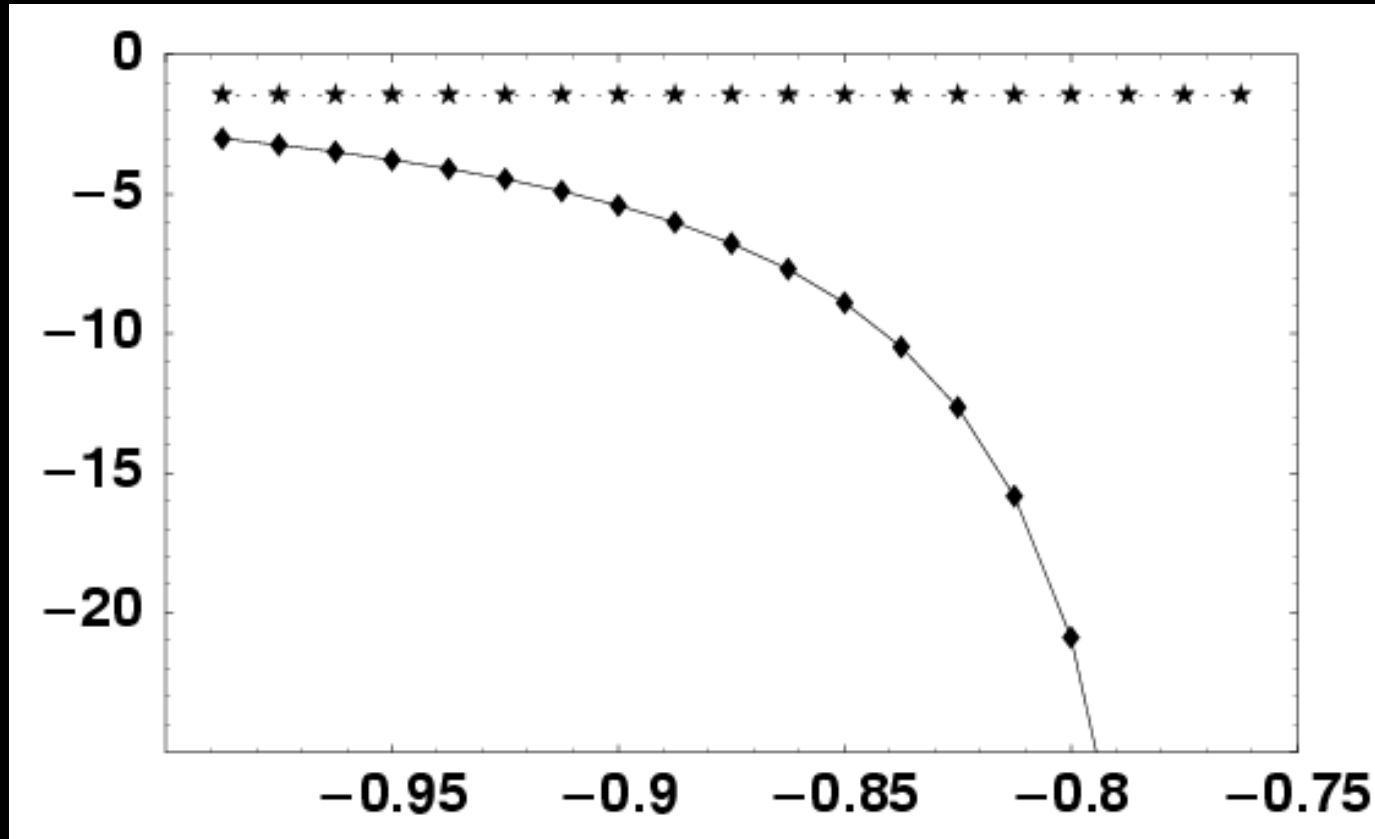


Use GSL!



The entropic Rip Stop!

$w \geq -1$ since today



If w switches to $w \geq -1$ today
The *GSL* still provides a constraint on w !



Summary and Conclusions

Conclusions



- In systems with gravitational interactions horizons contribute to the entropy.
 - A generalized second law of thermodynamics may hold.
 - The *GSL* is in conflict with big rip situations.
 - It places bounds on the equation of state
 - Viscosity is a possible but unlikely way out
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Viscosity

A way out?

$$S_{DE} \neq 0$$

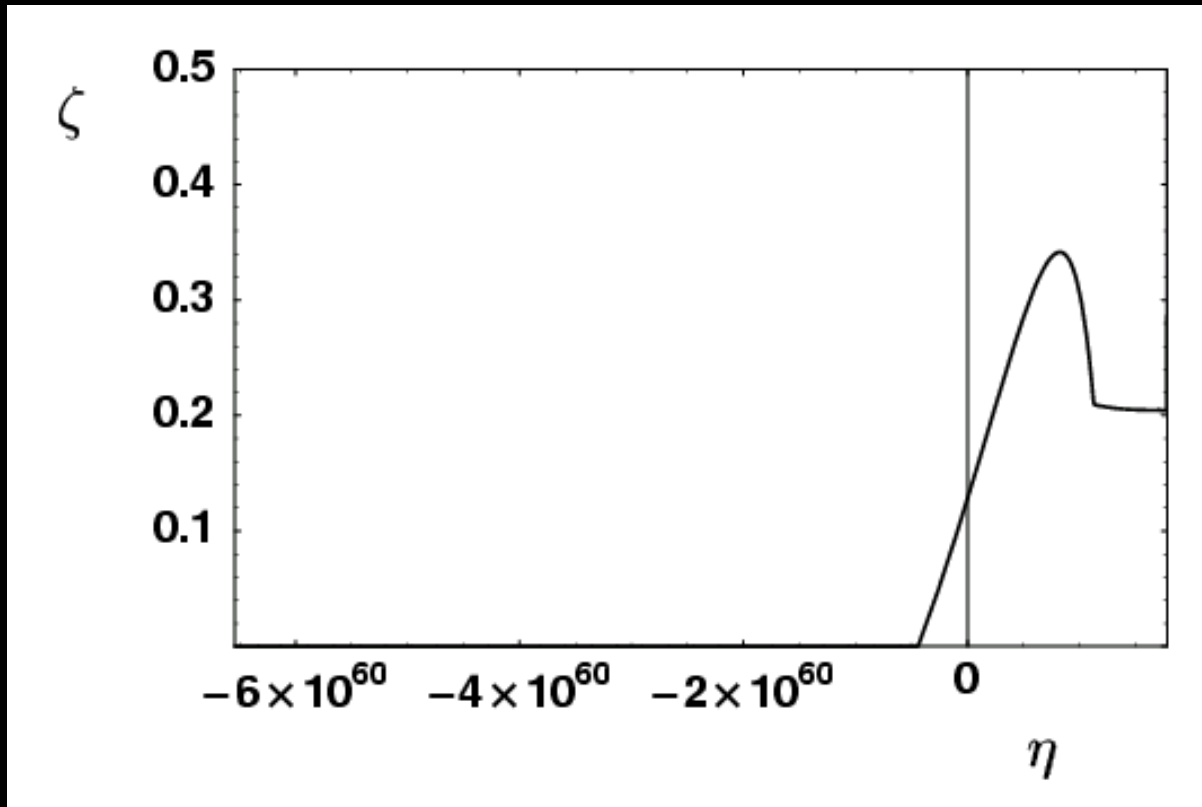


- So far we assumed that the dark energy is a scalar field or an adiabatically expanding perfect fluid.

- Bulk viscosity leads to entropy production!

$$\frac{1}{V} \frac{dS}{dt} = \frac{9H^2 \alpha \rho}{T} = \frac{3H\zeta\rho}{T}$$

Replenishing entropy



Natural η : 10^{-60}

Possible. But Huge viscosity is needed!